Modelling and Control of Nonlinear and Distributed Parameter Systems: The port-Hamiltonian Approach Inputs and Outputs

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February 26, 2020

Port-Hamiltonian systems with inputs and outputs

We are interested in boundary controls and boundary observations.

$$\begin{aligned} \frac{\partial x}{\partial t}(\zeta,t) &= \left(P_1 \frac{\partial}{\partial \zeta} + P_0\right) \left[\mathcal{H}x(t)\right] \\ \mathbf{u}(t) &= W_{B,1} \begin{bmatrix} (\mathcal{H}x)(b,t) \\ (\mathcal{H}x)(a,t) \end{bmatrix}, \\ 0 &= W_{B,2} \begin{bmatrix} (\mathcal{H}x)(b,t) \\ (\mathcal{H}x)(a,t) \end{bmatrix}, \\ \mathbf{y}(t) &= W_C \begin{bmatrix} (\mathcal{H}x)(b,t) \\ (\mathcal{H}x)(a,t) \end{bmatrix}. \end{aligned}$$

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$$\frac{\partial^2 w}{\partial t^2}(\zeta,t) = \frac{1}{\rho(\zeta)} \frac{\partial}{\partial \zeta} \left[T(\zeta) \frac{\partial w}{\partial \zeta}(\zeta,t) \right]$$

$$\frac{1}{\psi_y} \qquad u(t) = T(1) \frac{\partial w}{\partial \zeta}(1,t),$$

$$0 = \frac{\partial w}{\partial t}(0,t)$$

$$y(t) = \frac{\partial w}{\partial t}(1,t)$$

Question: Is this a well-posed linear system?

State space $X = L^2((a, b); \mathbb{R}^n)$ with (the energy) norm

$$\|f\|_X^2 = \frac{1}{2} \int_a^b f(\zeta)^T \mathcal{H}(\zeta) f(\zeta) d\zeta.$$

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Definition

The port-Hamiltonian system is called well-posed, if

•
$$Ax = P_1 \frac{d}{d\zeta} [\mathcal{H}x] + P_0 [\mathcal{H}x]$$
 with domain
 $D(A) = \left\{ x \in X \mid \frac{d}{d\zeta} \mathcal{H}x \in X, \begin{bmatrix} W_{B,1} \\ W_{B,2} \end{bmatrix} \begin{bmatrix} (\mathcal{H}x)(b) \\ (\mathcal{H}x)(a) \end{bmatrix} = 0 \right\}$
is the generator of a C_0 -semigroup on X .

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• There are
$$t_0, m_{t_0} > 0$$
:

$$\|x(t_0)\|_X^2 + \int_0^{t_0} \|y(t)\|^2 dt \le m_{t_0} \left[\|x(0)\|_X^2 + \int_0^{t_0} \|u(t)\|^2 dt \right]$$

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Theorem (Z, Le Gorrec, Maschke, Villegas '10) If $Ax = \left(P_1 \frac{d}{d\zeta} + P_0\right) [\mathcal{H}x]$ generates a C_0 -semigroup, then the port-Hamiltonian system is well-posed.

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Remark: We even have a regular system.



 $\bigvee_{y}^{u} \qquad \frac{\partial^{2} w}{\partial t^{2}}(\zeta, t) = \frac{1}{\rho(\zeta)} \frac{\partial}{\partial \zeta} \left[T(\zeta) \frac{\partial w}{\partial \zeta}(\zeta, t) \right]$ $u(t) = T(1) \frac{\partial w}{\partial \zeta}(1, t), \quad 0 = \frac{\partial w}{\partial t}(0, t)$ $y(t) = \frac{\partial w}{\partial t}(1,t)$ $P_1 \mathcal{H} = \begin{bmatrix} 0 & T \\ \frac{1}{a} & 0 \end{bmatrix} = \begin{bmatrix} \gamma & -\gamma \\ \frac{1}{a} & \frac{1}{a} \end{bmatrix} \begin{bmatrix} \gamma & 0 \\ 0 & -\gamma \end{bmatrix} \begin{vmatrix} \frac{1}{2\gamma} & \frac{\rho}{2} \\ -\frac{1}{2\gamma} & \frac{\rho}{2} \end{vmatrix} = S^{-1} \Delta S,$ with $\gamma > 0$ und $\gamma^2 = \frac{T}{a}$.

 $\begin{array}{ccc} & & & & \\ & & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ &$

with $\gamma > 0$ und $\gamma^2 = \frac{T}{\rho}$.

$$\begin{bmatrix} W_{B,1} \\ W_{B,2} \end{bmatrix} = \begin{bmatrix} 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \end{bmatrix}$$

 $\begin{array}{ccc}
 & & & \\$

$$P_1 \mathcal{H} = \begin{bmatrix} 0 & T \\ \frac{1}{\rho} & 0 \end{bmatrix} = \begin{bmatrix} \gamma & -\gamma \\ \frac{1}{\rho} & \frac{1}{\rho} \end{bmatrix} \begin{bmatrix} \gamma & 0 \\ 0 & -\gamma \end{bmatrix} \begin{bmatrix} \frac{1}{2\gamma} & \frac{\rho}{2} \\ -\frac{1}{2\gamma} & \frac{\rho}{2} \end{bmatrix} = S^{-1} \Delta S,$$

with $\gamma > 0$ und $\gamma^2 = \frac{T}{\rho}$.

$$\begin{bmatrix} W_{B,1} \\ W_{B,2} \end{bmatrix} = \begin{bmatrix} 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \end{bmatrix}$$

So if T and ρ are continuously differentiable, then the controlled wave equation is well-posed.

Inputs and Outputs Transfer Functions

Introduction

The aim of this part is to define transfer function for systems described by partial differential equations.

We derive these transfer functions via a very simple calculation. For port-Hamiltonian systems we show that the energy/power balance induces properties on the transfer function.

Transfer function for an o.d.e.

Consider the simple system described by the ordinary differential equation

$$\dot{y}(t) + 5y(t) = 3u(t),$$

the transfer function of this system is given by

$$G(s) = \frac{3}{s+5}.$$

How do you come to this?

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- Laplace transform, or
- Exponential solutions.

One way for obtaining the transfer function of

$$\dot{y}(t) + 5y(t) = 3u(t)$$

is to take $u(t) = e^{st}$, $s \in \mathbb{C}$, and to try to find a solution of the same format, i.e., $y(t) = \alpha e^{st}$.

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If $s \neq -5$, this is solvable;

$$\alpha = \frac{3}{s+5}.$$

So if we want to find an exponential solution

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- ▶ It is possible for all $s \in \mathbb{C}$ except for s = -5.
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▶ We call this the transfer function at *s*.

Definition

Given an (abstract) differential equation in the variables (u(t), z(t), y(t)), where u(t), z(t), and y(t) take their values in the (Hilbert) spaces U, Z, and Y, respectively. Let $s \in \mathbb{C}$. If for every $u_0 \in U$, there exists a unique solution of the form $(u_0e^{st}, z_0e^{st}, y_0e^{st})$, and the mapping $u_0 \mapsto y_0$ is linear and bounded, then this mapping is called the transfer function at s, and will be denoted by G(s).

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We call a solution of the form $(u_0e^{st}, z_0e^{st}, y_0e^{st})$ an exponential solution.

Transfer function for state linear systems

Consider the state differential equation

$$\dot{x}(t) = Ax(t) + Bu(t) y(t) = Cx(t) + Du(t)$$

with A, B, C, and D matrices.

Let $s \in \mathbb{C}$, and $u_0 \in U$. We try to find a solution of the form $(u(t), x(t), y(t)) = (u_0 e^{st}, x_0 e^{st}, y_0 e^{st}).$

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Since e^{st} is never zero, this is equivalent to:

$$(sI - A)z_0 = Bu_0$$

$$y_0 = Cz_0 + Du_0.$$

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If sI - A is invertible, then we find

$$y_0 = C(sI - A)^{-1}Bu_0 + Du_0.$$

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This clearly defines a bounded linear mapping from u_0 to y_0 , and so the transfer function at s is given by

$$G(s) = C(sI - A)^{-1}B + D.$$

This holds for all

 $s \in \rho(A) := \{s \in \mathbb{C} \mid (sI - A)^{-1} \text{ exists as bounded operator}\}.$

Example

We take a heated bar. We heat it uniformly at one half, and we measure (half) the average temperature in the other half;

$$\begin{aligned} \frac{\partial x}{\partial t}(\zeta,t) &= \frac{\partial^2 x}{\partial \zeta^2}(\zeta,t) + \mathbb{1}_{\left[\frac{1}{2},1\right]}(\zeta)u(t) \\ \frac{\partial x}{\partial \zeta}(0,t) &= \frac{\partial x}{\partial \zeta}(1,t) = 0 \\ y(t) &= \int_0^{\frac{1}{2}} x(\zeta,t)d\zeta. \end{aligned}$$

We obtain the transfer function.

We try to find an exponential solution of the p.d.e. This gives the following equations

$$sx_{0}(\zeta)e^{st} = \frac{d^{2}x_{0}}{d\zeta^{2}}(\zeta)e^{st} + \mathbb{1}_{[\frac{1}{2},1]}(\zeta)u_{0}e^{st}$$
$$\frac{dx_{0}}{d\zeta}(0)e^{st} = \frac{dx_{0}}{d\zeta}(1)e^{st} = 0$$
$$y_{0}e^{st} = \int_{0}^{\frac{1}{2}}x_{0}(\zeta)e^{st}d\zeta.$$

Hence

$$sx_{0}(\zeta) = \frac{d^{2}x_{0}}{d\zeta^{2}}(\zeta) + \mathbb{1}_{[\frac{1}{2},1]}(\zeta)u_{0}$$
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The first two lines represent an o.d.e. with boundary conditions.

The solution of

$$sx_{0}(\zeta) = \frac{d^{2}x_{0}}{d\zeta^{2}}(\zeta) + \mathbb{1}_{[\frac{1}{2},1]}(\zeta)u_{0}$$
$$\frac{dx_{0}}{d\zeta}(0) = \frac{dx_{0}}{d\zeta}(1) = 0$$

is given as

$$x_0(\zeta) = \cosh(\sqrt{s}\zeta)x_0(0) - \frac{1}{\sqrt{s}}\int_0^\zeta \sinh(\sqrt{s}(\zeta-\xi))\,\mathbb{1}_{[1/2,1]}(\xi)u_0d\xi$$

with

$$x_0(0) = \frac{\sinh(\sqrt{s}/2)u_0}{s\sinh(\sqrt{s})} = \frac{u_0}{2s\cosh(\sqrt{s}/2)}.$$

Using this we find that

$$y_0 = \int_0^{\frac{1}{2}} x_0(\zeta) d\zeta$$
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Hence the transfer function is given by

$$G(s) = \frac{\tanh(\sqrt{s}/2)}{2s\sqrt{s}}.$$

Transfer function, remark

If you write the solution of the o.d.e.

$$sx_{0}(\zeta) = \frac{d^{2}x_{0}}{d\zeta^{2}}(\zeta) + \mathbb{1}_{[\frac{1}{2},1]}(\zeta)u_{0}$$
$$\frac{dx_{0}}{d\zeta}(0) = \frac{dx_{0}}{d\zeta}(1) = 0$$

as a Fourier cosine series, then you find another expression for the transfer function. Namely,

$$G(s) = \frac{1}{4s} - 2\sum_{n=1}^{\infty} \frac{\sin(n\pi\frac{1}{2})^2}{n^2\pi^2(s+n^2\pi^2)}.$$

Transfer function, remark

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However, the transfer function is unique, and so we find that

$$\frac{\tanh(\sqrt{s}/2)}{2s\sqrt{s}} = G(s) = \frac{1}{4s} - 2\sum_{n=1}^{\infty} \frac{\sin(n\pi\frac{1}{2})^2}{n^2\pi^2(s+n^2\pi^2)}.$$

So we have seen that working with exponential solutions, directly on the p.d.e., works very well.

Note that it is (almost) the same as the engineering trick of replacing derivative with respect to time by an s.

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We can do that for systems with control and observation at the boundary.

Transfer function, boundary control and observation

Example

Consider the system with boundary control and observation

$$\begin{array}{rcl} \frac{\partial w}{\partial t}(\zeta,t) &=& \frac{\partial w}{\partial \zeta}(\zeta,t) \\ w(1,t) &=& u(t) \\ y(t) &=& w(0,t). \end{array}$$

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Substituting exponential functions for all signals, gives

$$sx_0(\zeta)e^{st} = \frac{dx_0}{d\zeta}(\zeta)e^{st}$$
$$x_0(1)e^{st} = u_0e^{st}$$
$$y_0e^{st} = x_0(0)e^{st}.$$

Thus

Example of transfer function with boundary control and observation

$$sx_0(\zeta) = \frac{dx_0}{d\zeta}(\zeta)$$
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This is an ordinary differential equation with given (end) condition, u_0 and unknown (initial) condition, y_0 .

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This is an ordinary differential equation with given (end) condition, u_0 and unknown (initial) condition, y_0 . The solution equals $x_0(\zeta) = e^{s(\zeta-1)}u_0$. Thus $y_0 = e^{-s}u_0$. The transfer function equals

$$G(s) = e^{-s} \qquad s \in \mathbb{C}.$$

Bode and Nyquist plots

Similar like for rational function, we can draw the Bode and Nyquist plot of general transfer functions

Bode and Nyquist plots

Similar like for rational function, we can draw the Bode and Nyquist plot of general transfer functions For instance the Bode magnitude plot of

$$G(s) = \frac{\tanh(\sqrt{s}/2)}{2s\sqrt{s}} - \frac{1}{4s}$$



Consider the port-Hamiltonian system with input and outputs

$$\begin{split} \frac{\partial x}{\partial t}(\zeta,t) &= \left(P_1 \frac{\partial}{\partial \zeta} + P_0\right) [\mathcal{H}x(t)] \\ u(t) &= W_{B,1} \begin{bmatrix} (\mathcal{H}x)(b) \\ (\mathcal{H}x)(a) \end{bmatrix}, \quad 0 = W_{B,2} \begin{bmatrix} (\mathcal{H}x)(b) \\ (\mathcal{H}x)(a) \end{bmatrix}, \\ y(t) &= W_C \begin{bmatrix} (\mathcal{H}x)(b) \\ (\mathcal{H}x)(a) \end{bmatrix} \end{split}$$

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y(t) = W_C \begin{bmatrix} (\mathcal{H}x)(b) \\ (\mathcal{H}x)(a) \end{bmatrix}$$

Assume that the energy balance can be expressed in the inputs and outputs. That is

$$\dot{H}(t) = \begin{bmatrix} u(t)^{\top}, y(t)^{\top} \end{bmatrix} Q \begin{bmatrix} u(t) \\ y(t) \end{bmatrix}$$

with Q a symmetric matrix.

Since exponential solutions are solutions, the power balance

$$\dot{H}(t) = \begin{bmatrix} u(t)^{\top}, y(t)^{\top} \end{bmatrix} Q \begin{bmatrix} u(t) \\ y(t) \end{bmatrix}$$

also holds for these.

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<u>Remark:</u> Since the s is the exponential solutions may be complex, we have to write the power balance for complex valued solutions. The (complex) power balance equals

$$\dot{H}(t) = \left[u(t)^*, y(t)^*\right] Q \begin{bmatrix} u(t) \\ y(t) \end{bmatrix}.$$

Hence for the exponential solution the power balance can be written as

$$\begin{split} \dot{H}(t) &= \left[u(t)^*, y(t)^* \right] Q \begin{bmatrix} u(t) \\ y(t) \end{bmatrix} \\ &= \left[u_0^* e^{\overline{s}t}, y_0^* e^{\overline{s}t}, \right] Q \begin{bmatrix} u_0 e^{st} \\ y_0 e^{st} \end{bmatrix} \end{split}$$

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Since

$$H(t) = ||x(t)||_X^2 = \langle x(t), x(t) \rangle_X,$$

we find for the exponential solution that

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Or equivalently:

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Example: Wave equation

$$\frac{\partial^2 w}{\partial t^2}(\zeta, t) = \frac{1}{\rho(\zeta)} \frac{\partial}{\partial \zeta} \left[T(\zeta) \frac{\partial w}{\partial \zeta}(\zeta, t) \right]$$

$$\underbrace{ \begin{array}{rcl} & u \\ & y \end{array}}^{u} & u(t) = T(1) \frac{\partial w}{\partial \zeta}(1, t), \quad 0 = \frac{\partial w}{\partial t}(0, t)$$

$$\underbrace{ \begin{array}{rcl} & y(t) \end{array}}_{v} & = \frac{\partial w}{\partial t}(1, t)$$

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$$\dot{H}(t) = u(t)y(t) = \begin{bmatrix} u(t)^*, y(t)^* \end{bmatrix} \begin{bmatrix} 0 & \frac{1}{2} \\ \frac{1}{2} & 0 \end{bmatrix} \begin{bmatrix} u(t) \\ y(t) \end{bmatrix}$$

Transfer function for the vibrating string system

From the general result we find

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Since $\|x_0\|_X^2 \geq 0$ and $|u_0|^2 > 0,$ we find that for $\mathsf{Re}(s) > 0$ there holds

 $\operatorname{Re}(G(s)) \ge 0$

Thus G is positive real.