Modelling and Control of Nonlinear and Distributed Parameter Systems: The port-Hamiltonian Approach Solutions, existence and properties

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$$= \begin{bmatrix} 0 & 1 \\ 1 & 0 \end{bmatrix} \frac{\partial}{\partial \zeta} \left(\underbrace{\begin{bmatrix} \frac{1}{\rho(\zeta)} & 0 \\ 0 & T(\zeta) \end{bmatrix}}_{=\mathcal{H}} x(\zeta, t) \right).$$

The wave equation, state differential equation

So with the variables $x_1 = \rho \frac{\partial w}{\partial t}$ and $x_2 = \frac{\partial w}{\partial \zeta}$ the p.d.e. becomes

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We have generalised this to our class of first order port-Hamiltonian equations.

Port-Hamiltonian partial differential equations

Our model class are p.d.e.'s of the form

$$\frac{\partial x}{\partial t}(\zeta,t) = \left(P_1 \frac{\partial}{\partial \zeta} + P_0\right) \left[\mathcal{H}x(\zeta,t)\right]$$

with

•
$$x(\zeta, t) \in \mathbb{R}^n$$
, $\zeta \in [a, b]$, $t \ge 0$

▶ P_1 is an invertible, symmetric real $n \times n$ -matrix,

- ▶ P_0 is a skew-symmetric real $n \times n$ -matrix,
- $\mathcal{H}(\zeta)$ is a symmetric, invertible $n \times n$ -matrix with $mI \leq \mathcal{H}(\zeta) \leq MI$ for some m, M > 0.

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The energy/Hamiltonian is defined as

$$H(t) = H(x(\cdot, t)) = \frac{1}{2} \int_{a}^{b} x(\zeta, t)^{T} \mathcal{H}(\zeta) x(\zeta, t) d\zeta.$$

Power balance

For the Port-Hamiltonian p.d.e. with energy/Hamiltonian

$$H(x(\cdot,t)) = \frac{1}{2} \int_{a}^{b} x(\zeta,t)^{T} \mathcal{H}(\zeta) x(\zeta,t) d\zeta,$$

it is not hard to show that along solutions; homework

$$\dot{H}(t) = \frac{dH}{dt}(x(\cdot,t)) = \frac{1}{2} \left[(\mathcal{H}x)^T \left(\zeta,t\right) P_1\left(\mathcal{H}x\right)\left(\zeta,t\right) \right]_a^b$$

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Thus the change of internal energy goes via the boundary of the spatial domain, i.e. power balance.

The wave equation, energy



The energy is given by

$$H(t) = \frac{1}{2} \int_0^1 \rho(\zeta) \left(\frac{\partial w}{\partial t}(\zeta, t)\right)^2 + T(\zeta) \left(\frac{\partial w}{\partial \zeta}(\zeta, t)\right)^2 d\zeta$$

with ρ is the mass density, and T is Young's modulus.

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$$\frac{1}{2}\int_0^1 \rho(\zeta) \left(\frac{\partial w_0}{\partial t}(\zeta)\right)^2 + T(\zeta) \left(\frac{\partial w_0}{\partial \zeta}(\zeta)\right)^2 d\zeta < \infty.$$

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This indicates that our states must be functions (of the spatial variable)

We chooce for the <u>state</u> variables as $x_1 := \rho \frac{\partial w}{\partial t}$ (the momentum), $x_2 := \frac{\partial w}{\partial \zeta}$ (the strain).

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$$\frac{1}{2} \int_0^1 \begin{bmatrix} x_1(\zeta,t) \\ x_2(\zeta,t) \end{bmatrix}^T \begin{bmatrix} \frac{1}{\rho(\zeta)} & 0 \\ 0 & T(\zeta) \end{bmatrix} \begin{bmatrix} x_1(\zeta,t) \\ x_2(\zeta,t) \end{bmatrix} d\zeta.$$

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Since $MI > \mathcal{H}(\zeta) > mI$, we see that finite energy condition implies that the state should satisfy for all $t \ge 0$:

$$\int_0^1 \left\| \begin{bmatrix} x_1(\zeta,t) \\ x_2(\zeta,t) \end{bmatrix} \right\|^2 d\zeta < \infty$$

The functions $[0,1] \ni \zeta \mapsto f(\zeta) \in \mathbb{R}^2$ which satisfy

$$\int_0^1 \|f(\zeta)\|^2 \, d\zeta < \infty$$

form the linear space $L^2((0,1); \mathbb{R}^2)$.

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However, the "energy" is still used to measure the size of x, i.e., the norm

$$\|f\|_X^2 = \frac{1}{2} \int_0^1 f(\zeta)^\top \mathcal{H}(\zeta) f(\zeta) d\zeta.$$

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$$||f||_X^2 = \frac{1}{2} \int_0^1 f(\zeta)^\top \mathcal{H}(\zeta) f(\zeta) d\zeta.$$

This norm is linked with the inner product

$$\langle f,g \rangle_X = \frac{1}{2} \int_0^1 f(\zeta)^\top \mathcal{H}(\zeta) g(\zeta) d\zeta$$

We see that $||f||_X^2 = \langle f, f \rangle_X$.

So based on the energy of our system, we have chosen our state space as $X=L^2((0,1);\mathbb{R}^2)$ with the inner product

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We have that $||f||_X^2$ is precisely the energy. So our state X is also called the energy space, i.e, the space consisting of all state/shapes/.... with finite energy. Note that we already rewrote the p.d.e. model of the vibrating string in our state variables.

$$\frac{\partial}{\partial t} \begin{bmatrix} x_1 \\ x_2 \end{bmatrix} (\zeta, t) = \underbrace{\begin{bmatrix} 0 & 1 \\ 1 & 0 \end{bmatrix}}_{=P_1} \frac{\partial}{\partial \zeta} \left(\underbrace{\begin{bmatrix} \frac{1}{\rho(\zeta)} & 0 \\ 0 & T(\zeta) \end{bmatrix}}_{=\mathcal{H}} x(\zeta, t) \right)$$

The wave equation, change of view point

Instead of seeing the state as a function of time and space, we see it as a function of time (which at each time depends on the spatial variable). So

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Thus (short hand) $x(t) = T(t)x_0$. What properties do we expect from the solution mapping T(t)?

We denote the state space by X. Thus our solution map

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Properties

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- ▶ $T(t)x_0 \in X$ for all $x_0 \in X$. Thus $||x(t)||_X < \infty$ whenever $||x_0||_X < \infty$. In particular, $||x(t)||_X \le m(t)||x_0||$ for some function m(t).

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- ▶ $T(t)x_0 \in X$ for all $x_0 \in X$. Thus $||x(t)||_X < \infty$ whenever $||x_0||_X < \infty$. In particular, $||x(t)||_X \le m(t)||x_0||$ for some function m(t).
- For all $x_0 \in X$ there holds

$$\lim_{t \downarrow 0} \|T(t)x_0 - x_0\|_X = 0 \quad \text{or} \quad \lim_{t \downarrow 0} T(t)x_0 = x_0.$$

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It tells that the solution becomes more and more the initial state when time get smaller and smaller.

We introduce some notation. $\mathcal{L}(X)$ denotes the set of linear and bounded operators from X to X. Thus if $Q \in \mathcal{L}(X)$, then

• $Q(\alpha x_0 + \beta \tilde{x}_0) = \alpha Q(x_0) + \beta Q(\tilde{x}_0)$, and

• there exists a $q \ge 0$ such that for all $x_0 \in X$,

 $||Q(x_0)|| \le q ||x_0||.$
Semigroup

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Definition

A strongly continuous semigroup (<u>C₀-semigroup</u>) is an operator valued function, $(T(t))_{t>0}$, from $[0,\infty)$ to $\mathcal{L}(X)$ which satisfies

▶
$$T(0) = I$$

▶ $T(t)T(s) = T(t+s), \quad t, s \in [0, \infty)$

For all $x_0 \in X$ there holds

$$\lim_{t \downarrow 0} T(t) x_0 = x_0.$$

Semigroup, example

Let a be a (complex or real) number, then

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is a C_0 -semigroup on $X = \mathbb{R}$. Let A be a (square) matrix, then

$$T(t) := e^{At}$$

is a C_0 -semigroup on the state space $X = \mathbb{R}^n$. Homework

Contraction semigroup

Definition The C_0 -semigroup $(T(t))_{t\geq 0}$ is contraction semigroup if $||T(t)x_0|| \leq ||x_0||$ for all $t \geq 0$ and for all $x_0 \in X$.

It is a unitary group if

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 $t \rightarrow$

Solution pH-system

We known that $e^{At}x_0$ is the solution of

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How about our port-Hamiltonian partial differential equation? For our port-Hamiltonian equation we have that the state is directly linked to the energy. So the p.d.e. must tell us what happens with the energy.

We have that

$$\dot{H}(t) = \frac{1}{2} \left[(\mathcal{H}x)^T \left(\zeta, t\right) P_1 \left(\mathcal{H}x\right) \left(\zeta, t\right) \right]_a^b$$

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So the boundary conditions must tell us what happens with this term.

Given our port-Hamiltonian partial differential equation

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We need to add boundary conditions to this p.d.e. That are conditions in $x(\overline{\zeta,t})$ for ζ equal to a or b.

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$$W_B \begin{bmatrix} \mathcal{H}(b)x(b,t) \\ \mathcal{H}(a)x(a,t) \end{bmatrix} = 0.$$

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Question: Which boundary conditions lead to unique solutions?

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Question: Which boundary conditions lead to unique solutions? We answer this question by using semigroup theory. However, we do it only for contraction semigroups.

Port-Hamiltonian p.d.e., existence of solutions

Theorem (Le Gorrec, Maschke & Zwart '05, Jacob & Zwart '11)

Assume the (standard) conditions on P_0 , P_1 and \mathcal{H} . Assume further that W_B is a $n \times 2n$ matrix of full rank.

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The solution map is a unitary C_0 -group (i.e. $||T(t)x_0|| = ||x_0||$, $\forall x_0, \forall t$) if and only if

$$\dot{H} = 0.$$



$$\begin{array}{lcl} \displaystyle \frac{\partial^2 w}{\partial t^2}(\zeta,t) & = & \displaystyle \frac{1}{\rho(\zeta)} \frac{\partial}{\partial \zeta} \left[T(\zeta) \frac{\partial w}{\partial \zeta}(\zeta,t) \right] \\ \displaystyle \frac{\partial w}{\partial t}(0,t) & = & \displaystyle T(1) \frac{\partial w}{\partial \zeta}(1,t) = 0 \end{array}$$

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$$= \begin{bmatrix} 0 & 1 & 0 & 0\\0 & 0 & 1 & 0 \end{bmatrix} \begin{bmatrix} \mathcal{H}(1)x(1,t)\\\mathcal{H}(0)x(0,t) \end{bmatrix}.$$

Now we check the conditions.

▶
$$P_1 = \begin{bmatrix} 0 & 1 \\ 1 & 0 \end{bmatrix}$$
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 for all ζ , then

$$\mathcal{H}(\zeta) = \begin{bmatrix} \rho(\zeta)^{-1} & 0\\ 0 & T(\zeta) \end{bmatrix} \text{ satisfy } mI_2 \le \mathcal{H}(\zeta) \le MI_2.$$

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$$\blacktriangleright \dot{H} = 0.$$

Thus the solution map is a unitary group on the energy space.

Summary

What we have introduced is:

- A general concept of state and state space.
- ► A general concept of the solution map.

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- A general concept of state and state space.
- A general concept of the solution map.

<u>Question</u>: How to formulate state/state space for a partial differential equation?
 We study an example first.

On the spatial domain [0,1] consider the p.d.e.

$$\begin{aligned} \frac{\partial w}{\partial t}(\zeta,t) &= \frac{\partial w}{\partial \zeta}(\zeta,t), \qquad \zeta \in [0,1], \ t \ge 0, \\ w(1,t) &= 0 \\ w(\zeta,0) &= w_0(\zeta) \qquad \text{(given)}. \end{aligned}$$

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As state x(t) we choose w at a time t.

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- So $x(t) = w(\cdot, t)$, or $(x(t))(\zeta) = w(\zeta, t)$.
- As state space we choose $L^2(0,1)$.
- If we now introduce $\dot{x}(t) = \frac{\partial w}{\partial t}(\cdot, t)$ and $Ax(t) = \frac{\partial w}{\partial \zeta}(\cdot, t)$, then the p.d.e. becomes

$$\dot{x}(t) = Ax(t).$$

State differential equation

So the p.d.e.

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can with $x(t) = w(\cdot, t)$, $\dot{x}(t) = \frac{\partial w}{\partial t}(\cdot, t)$, and $Ax(t) := \frac{\partial w}{\partial \zeta}(\cdot, t)$, be written as abstract differential equation:

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Where is the boundary condition?
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Where is the boundary condition? Another problem: The (spatial) derivative does not exist for all $x(t) \in L^2(0,1)$.

$\mathsf{More} \, \operatorname{on} \, A$

We see that A is a mapping working for a fixed t, i.e., so for $f\in L^2(0,1)$ we can define Af as

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More on A

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Summary on \boldsymbol{A}

So the p.d.e.

$$\begin{aligned} \frac{\partial w}{\partial t}(\zeta,t) &= \frac{\partial w}{\partial \zeta}(\zeta,t), \qquad \zeta \in [0,1], \ t \ge 0, \\ w(1,t) &= 0 \\ w(\zeta,0) &= w_0(\zeta) \end{aligned}$$

is written as the abstract differential equation:

$$\label{eq:constraint} \begin{split} \dot{x}(t) &= A x(t), \qquad x(0) = x_0 = w_0 \\ \text{with } x(t) &= w(\cdot,t) \in X = L^2(0,1) \text{, and} \\ (Af)\left(\zeta\right) &= \frac{df}{d\zeta}(\zeta) \end{split}$$

with domain:

$$D(A) = \{ f \in X \mid \frac{df}{d\zeta} \in X, f(1) = 0 \}.$$

Port-Hamiltonian p.d.e., state space

Given our port-Hamiltonian partial differential equation with boundary conditions

$$\frac{\partial x}{\partial t}(\zeta, t) = \left(P_1 \frac{\partial}{\partial \zeta} + P_0\right) \left[\mathcal{H}(\zeta) x(\zeta, t)\right]$$
$$0 = W_B \begin{bmatrix} \mathcal{H}(b) x(b, t) \\ \mathcal{H}(a) x(a, t) \end{bmatrix}$$

with the properties on P_0 , P_1 and \mathcal{H} .

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- As state we choose $x(t) = x(\cdot, t)$.
- ► As state space we choose the energy space, i.e., $X = L^2((0,1); \mathbb{R}^n)$ with inner product

$$\langle f,g\rangle_X = \frac{1}{2}\int_a^b f(\zeta)^\top \mathcal{H}(\zeta)g(\zeta)d\zeta.$$

Port-Hamiltonian p.d.e., state space formulation

With the state $x(t) = x(\cdot, t)$ and $X = L^2((0, 1); \mathbb{R}^n)$ our port-Hamiltonian p.d.e. with boundary conditions;

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$$\dot{x}(t) = Ax(t),$$

where

$$\mathbf{A}x = \left(P_1\frac{d}{d\zeta} + P_0\right)\left[\mathcal{H}x\right]$$

with domain

$$\frac{D(A)}{d\zeta} = \left\{ x \in X \mid \frac{d}{d\zeta}(\mathcal{H}x) \in X, W_B \begin{bmatrix} \mathcal{H}(b)x(b) \\ \mathcal{H}(a)x(a) \end{bmatrix} = 0 \right\}.$$

We have now written our p.d.e.'s as

$$\dot{x}(t) = Ax(t)$$

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What is their relation?

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What is their relation?

We begin by studying the question when X is finite-dimensional.

$\mathsf{Finding}\ A$

Let A be given as

$$A = \left(\begin{array}{cc} 1 & 2\\ 0 & 3 \end{array}\right),$$

then

$$e^{At} = \left(\begin{array}{cc} e^t & e^{3t} - e^t \\ 0 & e^{3t} \end{array} \right).$$

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<u>Problem:</u> Suppose now that you know only e^{At} . How would you find A back? <u>Answer</u> Evaluate the derivative of the semigroup at t = 0.

Since $\frac{d}{dt}e^{At} = Ae^{At}$, we have

$$\frac{d}{dt}e^{At}\mid_{t=0} = A.$$

Theorem

we have that

Assume that $(T(t))_{t\geq 0}$ is the solution map of our p.d.e., then for those $x_0 \in X$ for which the following limit exists

$$\lim_{t \downarrow 0} \frac{T(t)x_0 - x_0}{t},$$

$$Ax_0 = \lim_{t \to 0} \frac{T(t)x_0 - x_0}{t}$$

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Furthermore, D(A) consists of precisely those $x_0 \in X$ for which the limit exists. A is named the infinitesimal generator of the C_0 -semigroup

 $(T(t))_{t\geq 0}.$

Lemma

If $x_0 \in D(A)$, then for t > 0, $T(t)x_0$ is differentiable, and

$$\frac{d}{dt}\left(T(t)x_0\right) = AT(t)x_0.$$

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For $x_0 \in X$, $T(t)x_0$ is called a weak solution.

So given the (general) C_0 -semigroup $(T(t))_{t\geq 0}$, we could try to find A by differentiating it at t = 0.

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However, A is know (or can be defined from the p.d.e.). So the natural question is how to find $(T(t))_{t>0}$ from A.

Note there is a difference between knowing the existence of a solution and having the form/expression of the solution. The expression for the solution can be hard/impossible to find. So we concentrate on <u>existence</u>.

We do this for a special class of C_0 -semigroups.

Recall: Contraction semigroup

Definition

The C_0 -semigroup $(T(t))_{t\geq 0}$ is contraction semigroup if

 $||T(t)x_0|| \le ||x_0||$ for all $t \ge 0$ and for all $x_0 \in X$.

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We known that

$$||T(t)x_0||^2 = \langle T(t)x_0, T(t)x_0 \rangle.$$

For $x_0 \in D(A)$, we have that the derivative of $T(t)x_0$ equals $AT(t)x_0$. So if we differentiate $||T(t)x_0||^2$, we find

$$\frac{d}{dt} \|T(t)x_0\|^2 = \langle AT(t)x_0, T(t)x_0 \rangle + \langle T(t)x_0, AT(t)x_0 \rangle.$$

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Now we choose t = 0. We know that $T(0)x_0 = x_0$. Thus at time equal to zero, we find

$$\left. \frac{d}{dt} \left(\|T(t)x_0\|^2 \right) \right|_{t=0} = \langle Ax_0, x_0 \rangle + \langle x_0, Ax_0 \rangle.$$

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So if T(t) is a contraction semigroup, then

$$\langle Ax_0, x_0 \rangle + \langle x_0, Ax_0 \rangle = \frac{d}{dt} ||T(t)x_0||^2 |_{t=0} \le 0.$$

This has to hold for all $x_0 \in D(A)$.

Theorem (Lumer-Phillips)

Let A be a densely defined operator, then A generates a contraction semigroup on X if and only if

- 1. $\langle Ax_0, x_0 \rangle + \langle x_0, Ax_0 \rangle \leq 0$ for all $x_0 \in D(A)$.
- 2. The range of A I is the whole of X.

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Condition 1 comes from $\frac{d}{dt} ||T(t)x_0||^2 \leq 0$. So for pH this is equivalent to $\dot{H}(t) \leq 0$. Note that Condition 2 seems to be missing in our existence theorem for pH systems.

Example

Consider on the state space ${\cal X}=L^2(0,1)$ the operator ${\cal A}$ which is given as

$$Af = \frac{df}{d\zeta}, \qquad \zeta \in [0,1]$$

with the domain

$$D(A) = \left\{ f \in L^2(0,1) \mid f \text{ is absolutely continuous,} \\ \frac{df}{d\zeta} \in L^2(0,1) \text{ and } f(1) = 0 \right\}$$

Let us check the properties:
• A is densely defined in $L^2(0,1)$.

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 $\langle Ax, x \rangle + \langle x, Ax \rangle$

=

• A is densely defined in $L^2(0,1)$.

$$\begin{aligned} Ax, x\rangle + \langle x, Ax\rangle \\ &= \int_0^1 \frac{dx}{d\zeta}(\zeta) \overline{x(\zeta)} d\zeta + \int_0^1 x(\zeta) \overline{\frac{dx}{d\zeta}(\zeta)} d\zeta \end{aligned}$$

▶ A is densely defined in L²(0,1).

$$\langle Ax, x \rangle + \langle x, Ax \rangle$$

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▶ To see if the range of (A - I) is everything, we have for every $f \in L^2(0, 1)$ to solve (A - I)x = f.

Solving
$$(A - I)x = f$$
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$$x(\zeta) = -\int_{\zeta}^{1} e^{\zeta - \xi} f(\xi) d\xi.$$

Conclusion:

$$Af = \frac{df}{d\zeta}, \qquad \zeta \in [0,1]$$

with the domain

$$D(A) = \left\{ f \in L^2(0,1) \mid \frac{df}{d\zeta} \in L^2(0,1) \text{ and } f(1) = 0 \right\}$$

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generates a contraction semigroup on $X = L^2(0, 1)$. Note that A can also been seen as a pH system! Homework



Port-Hamiltonian Systems

Inputs and Outputs

Port-Hamiltonian systems with inputs and outputs

We are interested in boundary controls and boundary observations.

$$\frac{\partial x}{\partial t}(\zeta, t) = \left(P_1 \frac{\partial}{\partial \zeta} + P_0\right) [\mathcal{H}x(t)]$$
$$\boldsymbol{u}(t) = W_{B,1} \begin{bmatrix} (\mathcal{H}x)(b) \\ (\mathcal{H}x)(a) \end{bmatrix}, 0 = W_{B,2} \begin{bmatrix} (\mathcal{H}x)(b) \\ (\mathcal{H}x)(a) \end{bmatrix}, \boldsymbol{y}(t) = W_C \begin{bmatrix} (\mathcal{H}x)(b) \\ (\mathcal{H}x)(a) \end{bmatrix}$$

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Example: Wave equation

$$\frac{\partial^2 w}{\partial t^2}(\zeta, t) = \frac{1}{\rho(\zeta)} \frac{\partial}{\partial \zeta} \left[T(\zeta) \frac{\partial w}{\partial \zeta}(\zeta, t) \right]$$
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Question: Is this a well-posed linear system?

State space $X\!=\!L^2((a,b);\mathbb{R}^n)$ with (the energy) norm

$$\|f\|_X^2 = \frac{1}{2} \int_a^b f(\zeta)^T \mathcal{H}(\zeta) f(\zeta) d\zeta.$$

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The port-Hamiltonian system is called well-posed, if

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$$Ax = P_1 \frac{d}{d\zeta} [\mathcal{H}x] + P_0 [\mathcal{H}x]$$
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• There are
$$t_0, m_{t_0} > 0$$
:

$$\|x(t_0)\|_X^2 + \int_0^{t_0} \|y(t)\|^2 dt \le m_{t_0} \left[\|x(0)\|_X^2 + \int_0^{t_0} \|u(t)\|^2 dt \right]$$

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Theorem (Z, Le Gorrec, Maschke, Villegas '10) If $Ax = \left(P_1 \frac{d}{d\zeta} + P_0\right) [\mathcal{H}x]$ generates a C_0 -semigroup, then the port-Hamiltonian system is well-posed.

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Remark: We even have a regular system.



 $\bigvee_{y}^{u} \qquad \frac{\partial^{2} w}{\partial t^{2}}(\zeta, t) = \frac{1}{\rho(\zeta)} \frac{\partial}{\partial \zeta} \left[T(\zeta) \frac{\partial w}{\partial \zeta}(\zeta, t) \right]$ $u(t) = T(1) \frac{\partial w}{\partial \zeta}(1, t), \quad 0 = \frac{\partial w}{\partial t}(0, t)$ $y(t) = \frac{\partial w}{\partial t}(1,t)$ $P_1 \mathcal{H} = \begin{bmatrix} 0 & T \\ \frac{1}{a} & 0 \end{bmatrix} = \begin{bmatrix} \gamma & -\gamma \\ \frac{1}{a} & \frac{1}{a} \end{bmatrix} \begin{bmatrix} \gamma & 0 \\ 0 & -\gamma \end{bmatrix} \begin{vmatrix} \frac{1}{2\gamma} & \frac{\rho}{2} \\ -\frac{1}{2\gamma} & \frac{\rho}{2} \end{vmatrix} = S^{-1} \Delta S,$ with $\gamma > 0$ und $\gamma^2 = \frac{T}{a}$.

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 $\begin{array}{ccc}
 & & & \\$

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So if T and ρ are continuously differentiable, then the controlled wave equation is well-posed.

Exercises

1. Show that for a pH system there holds:

$$\dot{H}(t) = \frac{dH}{dt}(x(\cdot,t)) = \frac{1}{2} \left[\left(\mathcal{H}x\right)^T \left(\zeta,t\right) P_1\left(\mathcal{H}x\right)\left(\zeta,t\right) \right]_a^b$$

- 2. Show that e^{At} is a C_0 -semigroup, when A is a (square) matrix.
- 3. Show that $Af = \frac{df}{d\zeta}$ with domain $D(A) = \{f \in L^2(0,1) \mid f \text{ is such that } \frac{df}{d\zeta} \in L^2(0,1) \text{ and } f(1) = 0\}$ can be associated to a pH system.

Exercise

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- a Show that the connected wave equations shown below can be written as a pH system,
 - b Show that for no force (u = 0) we have that the solution map is a contraction semigroup.
 - c Assume that we measure the velocity of the (vertical moving) middle bar. Show that the system is well-posed.

