

# Computing Normal Forms of quadratic DAE

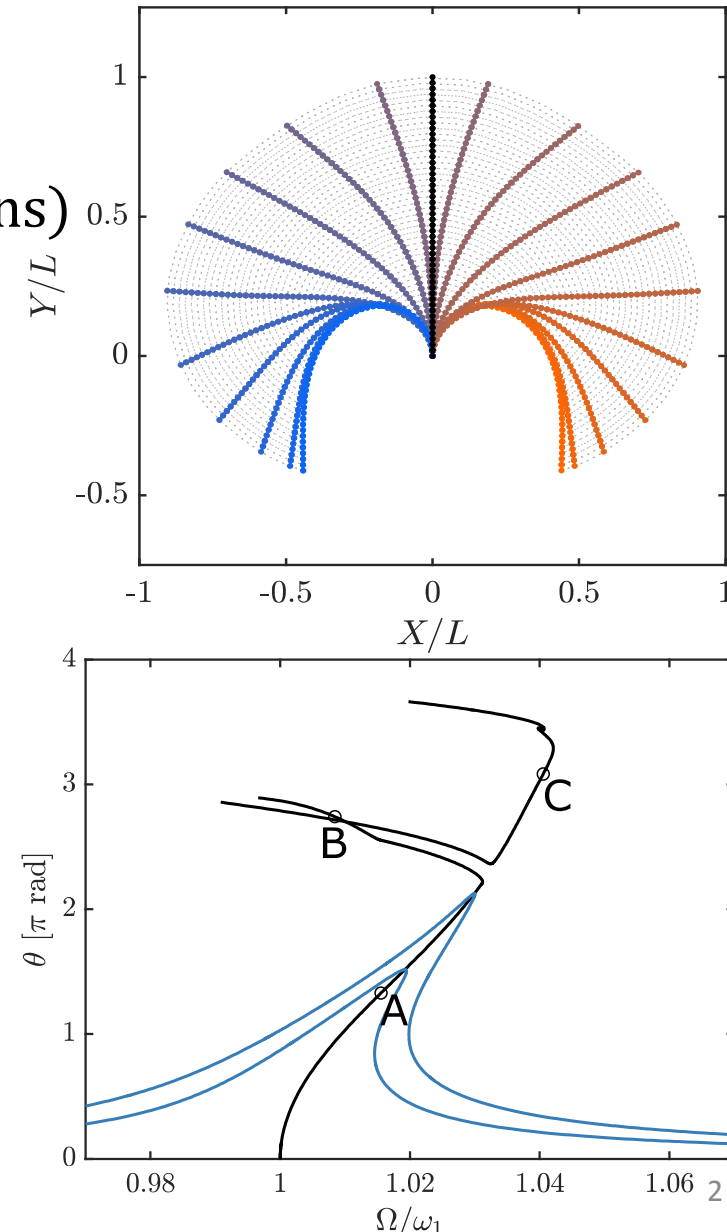
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# Motivation

- Work on geometrically exact beam FEM
  - M. Debeurre PhD (Thread project)
  - Use of constraints to parameterize rotations (unit quaternions)
  - Model reduction
- Discussion with A.Vizzacaro:
  - Normal form theory on ODE (with C. Touzé)
  - Can we compute normal form of constrained system (DAE) ?
- Work with MANLAB (B.Cochelin, L.Guillot)
  - Computation and continuation of periodic solutions
  - Quadratic DAE as an input
- Can we compute normal form of quadratic DAE ?



# Quadratic Differential Algebraic Equation

- Quadratic DAE (MANLAB input)

$$A\dot{y} = Ly + Q(y, y)$$

- $y$  is the vector of unknowns (N variables)
  - $A$  is the « mass » operator (NxN constant matrix, possibly singular)
  - $L$  is the « stiffness » operator (NxN constant matrix)
  - $Q$  is a bilinear operator (quadratic)
- A very general way of writing non linear dynamical equation:
    - Deals with constrained system (Lagrange multipliers)
    - Deals with any kind of nonlinearity, provided one finds a quadratic recast

# Normal Form Theory

- Original dynamics (quadratic DAE)

$$A\dot{y} = Ly + Q(y, y)$$

- Normal form:

- Find a polynomial change of variable

$$y = W(z) = P_d(z) + o(z^d)$$

- Find the polynomial normal dynamics under simplest form (up to order  $d$ ):

$$\dot{z} = f(z) + o(z^d)$$

- Such that the original dynamics is exact (up to order  $d$ )

- In practice:

- Derivatives (chain rule):

$$\dot{y} = \frac{\partial W}{\partial z} \dot{z} = \frac{\partial W}{\partial z} f$$

- Homological equation:

$$A \frac{\partial W}{\partial z} f = LW + Q(W, W) + o(z^d)$$

- Allows to compute  $W$  and  $f$

# Normal Form Theory

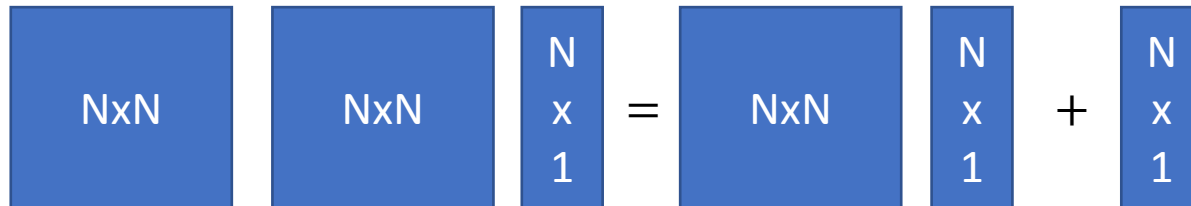
- Full normal form:

- Change of variable:

$$y = W(z)$$

- $z$  contains as many variables as  $y$
- $N$  variables in total
- Homological equation:

$$A \frac{\partial W}{\partial z} f = LW + Q(W, W)$$



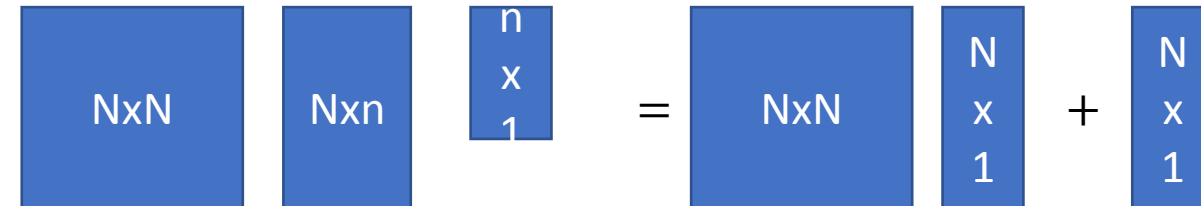
- Reduced normal form:

- Change of variable:

$$y = W(z)$$

- $z$  is much smaller in size than  $y$
- $n \ll N$  (model reduction)
- Homological equation:

$$A \frac{\partial W}{\partial z} f = LW + Q(W, W)$$



# Polynomials representation and normal form

- Polynomial of deg.  $d$  in  $n$  variables:
  - Vector space of dimension  $M$
  - Monomial Basis  $B = B(z) = (b_1, \dots, b_M)$
  - Vectors of polynomials

- Change of coordinates:

$$y = W(z) = \sum_{1 \leq m \leq M} W_m b_m(z),$$

Coefficient :  $W_m \in \mathbb{C}^N$

- Reduced normal dynamics :

$$\dot{z} = f(z) = \sum_{1 \leq m \leq M} f_m b_m(z),$$

Coefficient :  $f_m \in \mathbb{C}^n$

- Homological equation:

$$A \frac{\partial W}{\partial z} f = LW + Q(W, W)$$

- Substitute polynomial expression and balance monomials coefficients

- Needs multiplication and derivative:
  - Linear operations
  - Use matrix representation (need the effect of multiplication and derivative on the basis elements)

# Polynomials: Derivative and multiplication

- Consider the space of polynomial of degree  $d$  in  $n$  variables  $z = (z_1, \dots, z_n)$ :
  - Vector space of dimension  $M$
  - Basis  $B = B(z) = (b_1, \dots, b_M)$

- Derivatives:

- There exists  $n$  « derivative matrices »  $\nabla_k$  such that:

$$\frac{\partial B^T}{\partial z_k} = \nabla_k B^T$$

- Multiplication:

- There exists  $M$  « multiplication matrices »  $\Lambda_r$  such that:

$$b_r B^T = \Lambda_r B^T$$

- Advantages of the vector space structure:

- Multiplication and derivation matrices are sparse and can be computed prior to the normal form computation
  - Use linear algebra to write explicitly the normal form homological equation

# Polynomials: application to normal form

- Polynomial of deg.  $d$  in  $n$  variables:

- Basis  $B = B(z) = (b_1, \dots, b_M)$

- Homological equation:

$$A \frac{\partial W}{\partial z} f = LW + Q(W, W)$$

- Change of coordinates:

$$y = W(z) = \sum_{1 \leq m \leq M} W_m b_m(z),$$

Coefficient :  $W_m \in \mathbb{C}^N$

- Substitute polynomials expression :

$$AW \sum_{m=1}^{M-1} \sum_{j=1}^n f_m^j \nabla^j \Delta^m = LW + \sum_{m=1}^{M-1} \sum_{s=1}^{M-1} Q(W_m, W_s) b_m \Delta^s$$

- Reduced normal dynamics :

$$\dot{z} = f(z) = \sum_{1 \leq m \leq M} f_m b_m(z),$$

Coefficient :  $f_m \in \mathbb{C}^n$

- Equation for coeff. of monomial  $b_k$ :

$$(\sigma_{kk} A - L) W_k + \sum_{m \neq k} \sigma_{km} A W_m = \sum_{m=1}^M \sum_{l=1}^M \Lambda_{lk}^m Q(W_m, W_l)$$



# Solution to the homological equation

- Sequential resolution of the homological equation for each degree
- Order 0 (constant monomials:  $b_1 = 1$ ):
  - Static solution
- Order 1 (linear monomials :  $b_2 = z_1 \dots, b_{n+1} = z_n$ ):
  - Linear modes of interest (complex mode shapes  $Y_k$ )
  - $W_2 = Y_k, \dots, W_{n+1} = Y_k$
- Higher Orders:

$$(\sigma_{kk}\mathbf{A} - \mathbf{L}) \mathbf{W}_k + \sum_{m \neq k} \sigma_{km} \mathbf{A} \mathbf{W}_m = \sum_{m=1}^M \sum_{l=1}^M \Lambda_{lk}^m \mathbf{Q}(\mathbf{W}_m, \mathbf{W}_l)$$

- After close inspection of the terms:

$$\begin{bmatrix} \sigma_{kk}\mathbf{A} - \mathbf{L} & \mathbf{A} \mathbf{Y}_{r \in \mathcal{R}(k)} & \mathbf{0} \\ \mathbf{X}_{r \in \mathcal{R}(k)}^T \mathbf{A} & \mathbf{0} & \mathbf{0} \\ \mathbf{0} & \mathbf{0} & \mathbf{I} \end{bmatrix} \begin{bmatrix} \mathbf{W}_k \\ \mathbf{f}_{k, s \in \mathcal{R}(k)} \\ \mathbf{f}_{k, s \notin \mathcal{R}(k)} \end{bmatrix} = \begin{bmatrix} \mathbf{RHS}_k \\ \mathbf{0} \\ \mathbf{0} \end{bmatrix}$$

# Normal form of DAE sum up

- Original dynamics (quadratic DAE)

$$A\dot{y} = Ly + Q(y, y)$$

- Normal form:

- change of variable

$$y = W(z) = \sum_{1 \leq m \leq M} W_m b_m(z),$$

- normal dynamics under simplest form :

$$\dot{z} = f(z) = \sum_{1 \leq m \leq M} f_m b_m(z)$$

- Compute derivation ( $\nabla_r$ ) and multiplication ( $\Lambda_s$ ) matrices

- Homological equation (monom.  $b_k$ ):

$$\begin{bmatrix} \sigma_{kk} A - L & AY_{r \in \mathcal{R}(k)} & 0 \\ X_{r \in \mathcal{R}(k)}^T A & 0 & 0 \\ 0 & 0 & I \end{bmatrix} \begin{bmatrix} W_k \\ f_{k, s \in \mathcal{R}(k)} \\ f_{k, s \notin \mathcal{R}(k)} \end{bmatrix} = \begin{bmatrix} RHS_k \\ 0 \\ 0 \end{bmatrix}$$

- Solve for  $W_k$  and  $f_k$  sequentially

- Use reduced dynamics to compute solutions,

- Go back to physical variables using the change of variables

# Example 1: Duffing

- Duffing Oscillator

$$\ddot{u} + u + u^3 = 0$$

- Quadratic Equation of motion

$$\begin{pmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 0 \end{pmatrix} \begin{pmatrix} \dot{u} \\ \dot{v} \\ \dot{r} \end{pmatrix} = \begin{pmatrix} 0 & 1 & 0 \\ -1 & 0 & 0 \\ 0 & 0 & 1 \end{pmatrix} \begin{pmatrix} u \\ v \\ r \end{pmatrix} + \begin{pmatrix} 0 \\ -ur \\ -u^2 \end{pmatrix}$$

$$A\dot{y} = Ly + Q(y, y)$$

# Example 1: Duffing

- Linear modes

$$Y_1 = \begin{pmatrix} -i \\ 1 \\ 0 \end{pmatrix}, Y_2 = \begin{pmatrix} i \\ 1 \\ 0 \end{pmatrix}, Y_3 = \begin{pmatrix} 0 \\ 0 \\ 1 \end{pmatrix} \quad \lambda_1 = i, \lambda_2 = -i, \lambda_3 = \infty$$

- Normal form (ordre 1)

- Linear change of variable

$$W_1 = Y_1 \text{ and } W_2 = Y_2 \text{ and } f_{11} = i, f_{22} = -i \text{ and } f_{12} = f_{21} = 0.$$

- Uncoupled linear dynamics

$$\dot{z} = \begin{pmatrix} \dot{z}_1 \\ \dot{z}_2 \end{pmatrix} = f(z) = \begin{pmatrix} iz_1 \\ -iz_2 \end{pmatrix}$$

# Example 1: Duffing

- Normal form (ordre 2)

- Coeff  $z_1^2$  ( $2\lambda_1 = 2i$ ) :  $f_3 = 0$

$$W_3 = (2iA - L)^{-1} Q(Y_1, Y_1)$$

$$W_3 = \begin{pmatrix} 0 \\ 0 \\ -1 \end{pmatrix}$$

- Coeff  $z_1 z_2$  ( $\lambda_1 + \lambda_2 = 0$ ) :  $f_4 = 0$

$$W_4 = -L^{-1} [Q(Y_1, Y_2) + Q(Y_2, Y_1)]$$

$$W_4 = \begin{pmatrix} 0 \\ 0 \\ 2 \end{pmatrix}$$

- Coeff  $z_2^2$  ( $2\lambda_2 = -2i$ ) :  $f_5 = 0$

$$W_5 = (-2iA - L)^{-1} Q(Y_2, Y_2)$$

$$W_5 = \begin{pmatrix} 0 \\ 0 \\ -1 \end{pmatrix}$$

At this point the reduced dynamics is still linear !

# Example 1: Duffing

- Normal form (ordre 3, non resonant terms)

- Coeff  $z_1^3$  ( $3\lambda_1 = 3i$ ) :  $f_6 = 0$

$$W_6 = (3iA - L)^{-1}(Q(Y_1, W_3) + Q(W_3, Y_1))$$

$$W_6 = \begin{pmatrix} i \\ \frac{1}{4} \\ \frac{3}{4} \\ \frac{1}{4} \\ 0 \end{pmatrix}$$

- Coeff  $z_2^3$  ( $3\lambda_2 = -3i$ ) :  $f_9 = 0$

$$W_9 = (-3iA - L)^{-1}(Q(Y_2, W_5) + Q(W_5, Y_2))$$

$$W_9 = \begin{pmatrix} i \\ \frac{1}{4} \\ \frac{3}{4} \\ \frac{1}{4} \\ 0 \end{pmatrix}$$

# Example 1: Duffing

- Normal form (ordre 3, resonant terms)

- Coeff  $z_1^2 z_2$  ( $2\lambda_1 + \lambda_2 = \lambda_1 = i$ )

$$(iA - L)W_7 + f_{17}AY_1 = Q(Y_1, W_4) + Q(W_3, Y_2) + Q(W_4, Y_1) + Q(Y_2, W_3)$$

$$X_1^T W_7 = 0$$

$$W_7 = \begin{pmatrix} \frac{3i}{4} \\ \frac{3}{4} \\ \frac{3}{4} \\ 0 \end{pmatrix}$$

- Coeff  $z_1 z_2^2$  ( $\lambda_1 + 2\lambda_2 = \lambda_2 = -i$ )

$$(-iA - L)W_8 + f_{28}AY_2 = Q(Y_1, W_5) + Q(W_5, Y_1) + Q(W_4, Y_2) + Q(Y_2, W_4)$$

$$X_2^T W_8 = 0$$

$$W_8 = \begin{pmatrix} -\frac{3i}{4} \\ \frac{3}{4} \\ \frac{3}{4} \\ 0 \end{pmatrix}$$

-> Addition of 2 nonlinear terms in the reduced dynamics

# Example 1: Duffing

- Normal form : Sum up
  - Change of coordiantes

$$y = \begin{pmatrix} u \\ v \\ r \end{pmatrix} = W(z) = \begin{pmatrix} i(-z_1 + z_2) + \frac{i}{4}z_1^3 + \frac{3i}{4}z_1^2z_2 - \frac{3i}{4}z_1z_2^2 + \frac{i}{4}z_2^3 \\ z_1 + z_2 + \frac{3}{4}z_1^3 + \frac{3}{4}z_1^2z_2 + \frac{3}{4}z_1z_2^2 + \frac{3}{4}z_2^3 \\ -z_1^2 + 2z_1z_2 + z_2^2 \end{pmatrix}$$

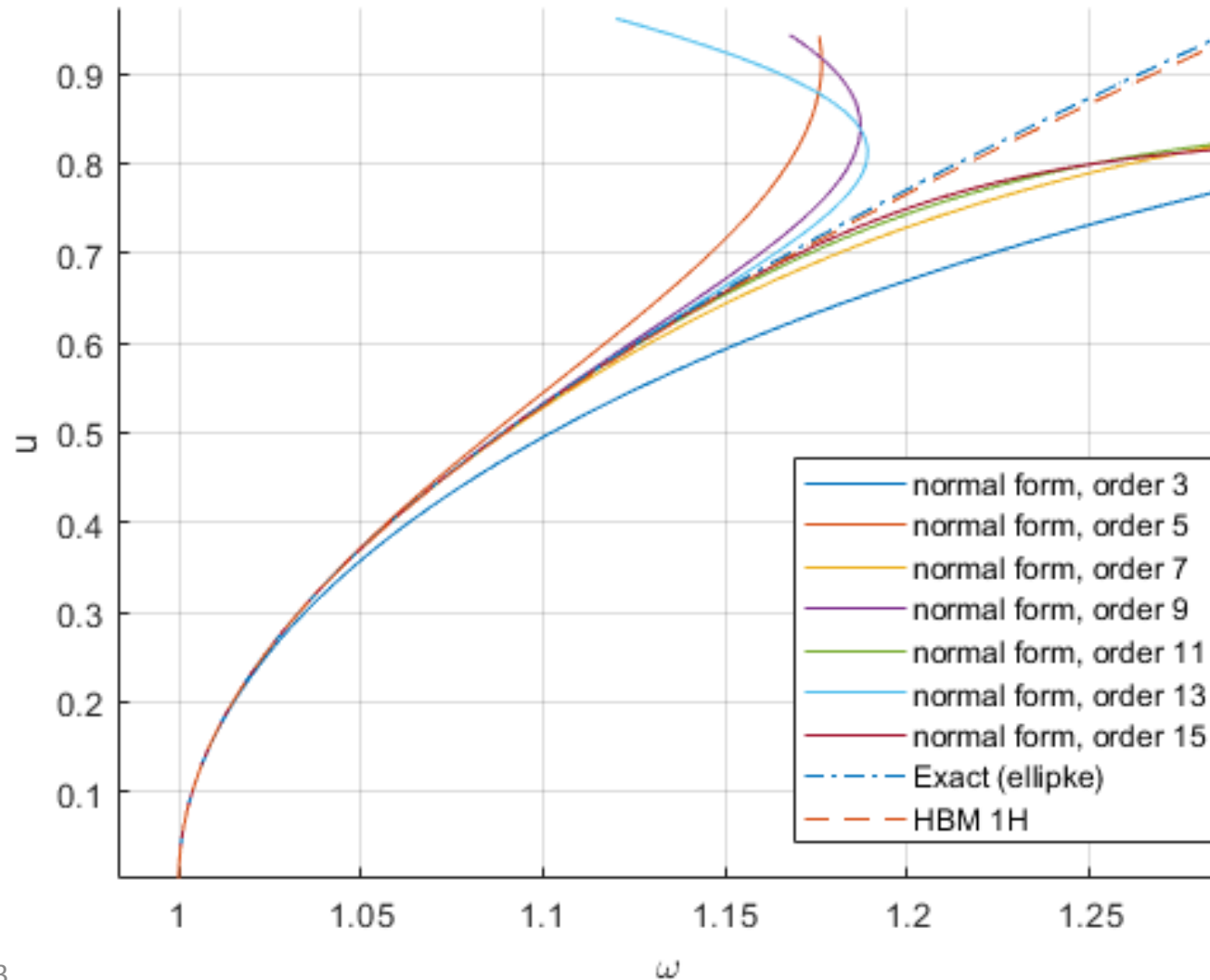
- Reduced dynamics:

$$\dot{z} = \begin{pmatrix} \dot{z}_1 \\ \dot{z}_2 \end{pmatrix} = f(z) = \begin{pmatrix} iz_1 + \frac{3i}{4}z_1^2z_2 \\ -iz_2 - \frac{3i}{4}z_1z_2^2 \end{pmatrix} \quad z_1 = \rho e^{i\theta} \quad \longrightarrow \quad \begin{cases} \dot{\rho} = 0 + o(\rho^3) \\ \dot{\theta} = 1 + \frac{3}{2}\rho^2 + o(\rho^3) \end{cases}$$



# Example 1: Duffing

- Backbone curve :



$$\dot{\rho} = 0 + o(\rho^3)$$
$$\dot{\theta} = 1 + \frac{3}{2}\rho^2 + o(\rho^3)$$

# Example 2: simple pendulum

- Simple pendulum (Unit length, mass and gravity)

- Parametrization ( $n \in \mathbb{N}$ ):

- $p_0 = \cos \frac{\theta}{n}, p_3 = \sin \frac{\theta}{n}$

- Kinetic Energy:

- $T = \frac{n^2}{2} (\dot{p}_0^2 + \dot{p}_3^2)$

- Potential Energy:

- $U = -T_n(p_0)$

- Constraint:

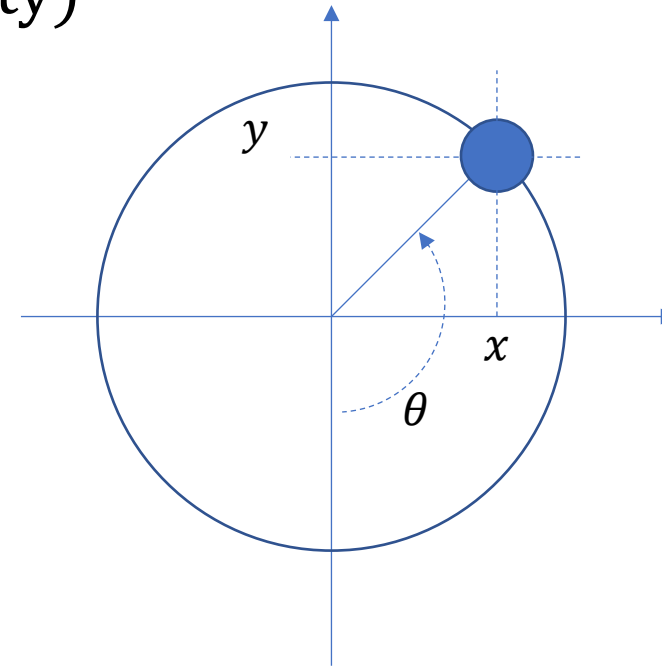
- $p_0^2 + p_3^2 - 1 = 0$

- Equation of motion

$$n^2 \ddot{p}_0 - \frac{\partial T_n}{\partial p_0} = 2\lambda p_0$$

$$n^2 \ddot{p}_3 = 2\lambda p_3$$

$$p_0^2 + p_3^2 - 1 = 0$$



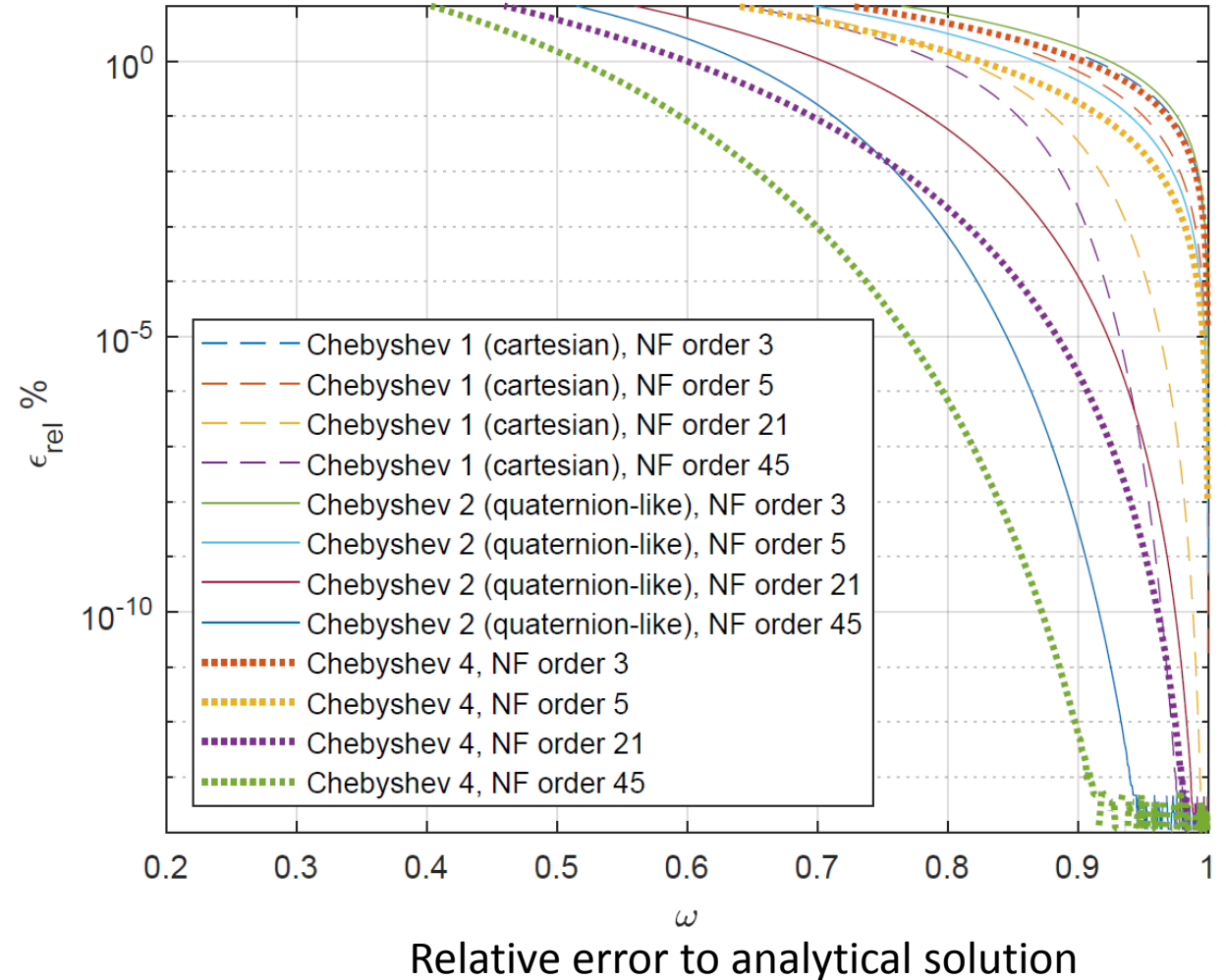
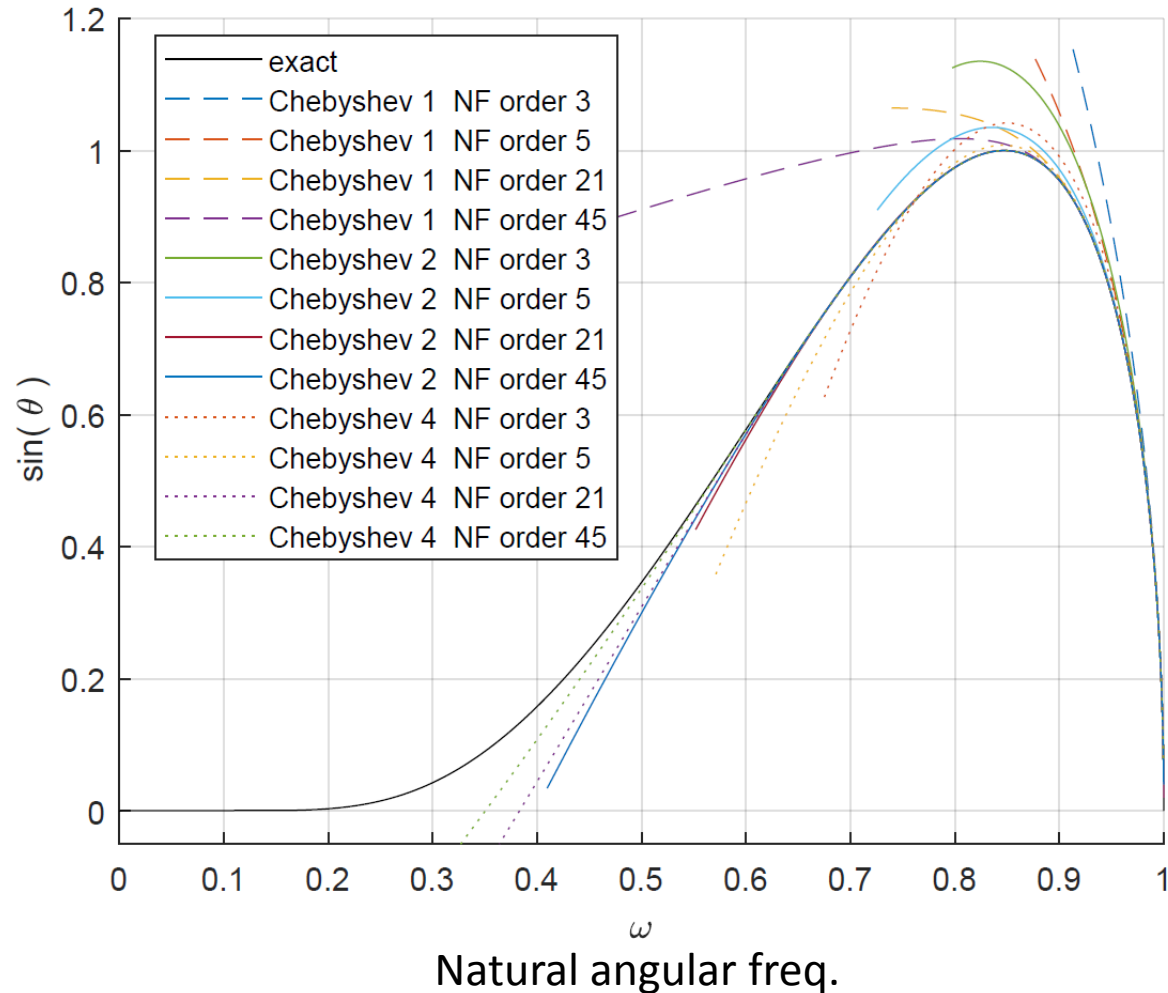
- $T_n$ :  $n$  – th Chebichev polynomials

- $n = 1$  :catesian parametrisation

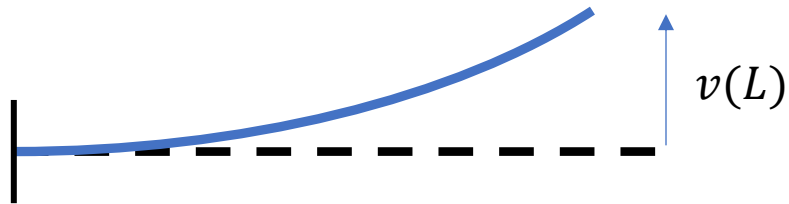
- $n = 2$  : quaternion like parametrisation

# Example 2: simple pendulum

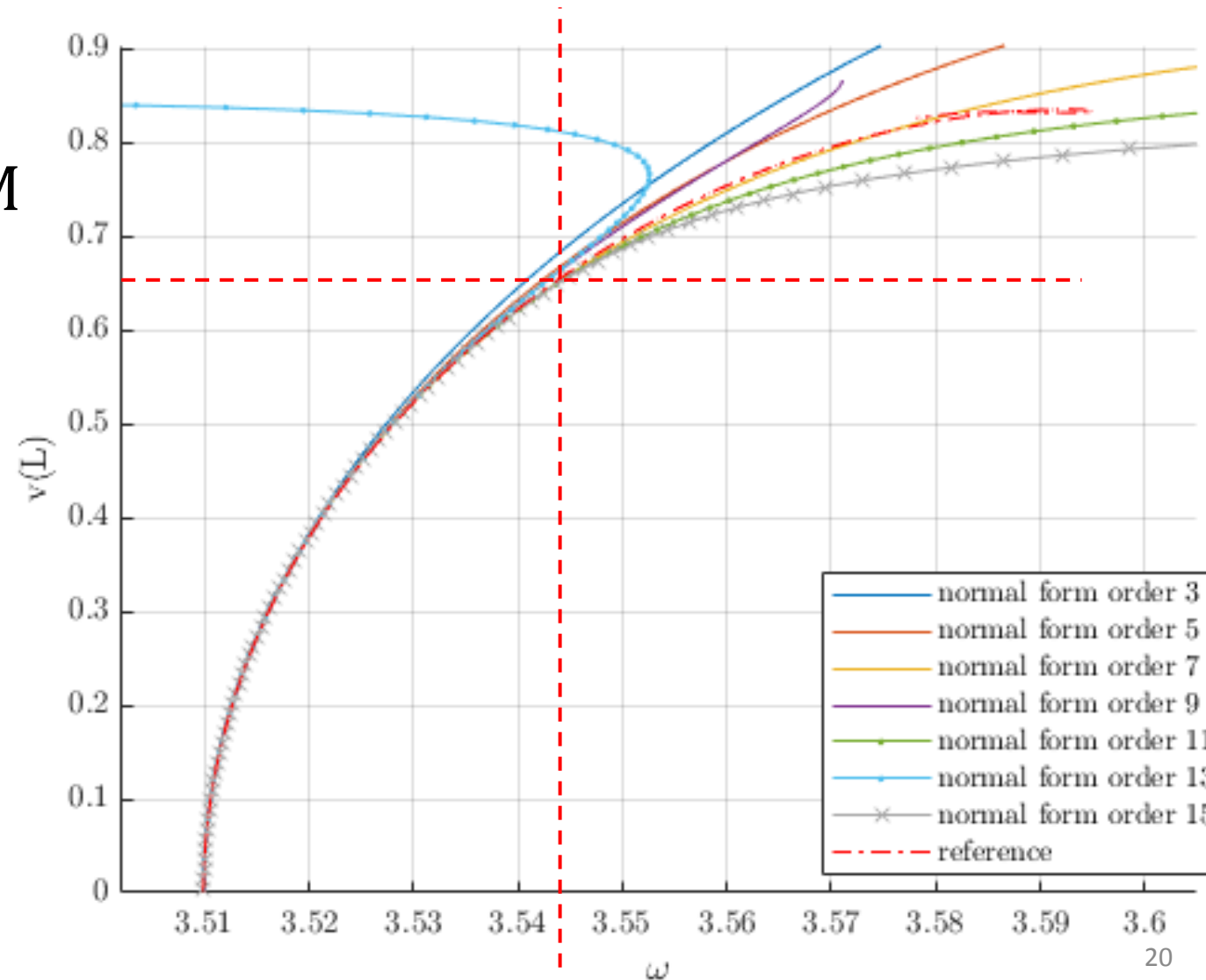
## • Results



# Example 3: geometrically exact beam FE



- 2D geometrically exact beam FEM
  - Parametrisation of rotation using « quaternions »
  - Rewritten under a quadratic DAE
- Cantilever straight beam
  - FEM: 30 quadratic elements
  - DAE:  $\sim 500$  variables
- Reference solution:
  - HBM: 20 Harmonics
- Normal form
  - 2 variables



# Quick word on Validity range

- Backbone

$$\omega(\rho) = \dot{\theta} = \sum a_n \rho^n$$

- Radius of convergence of a series

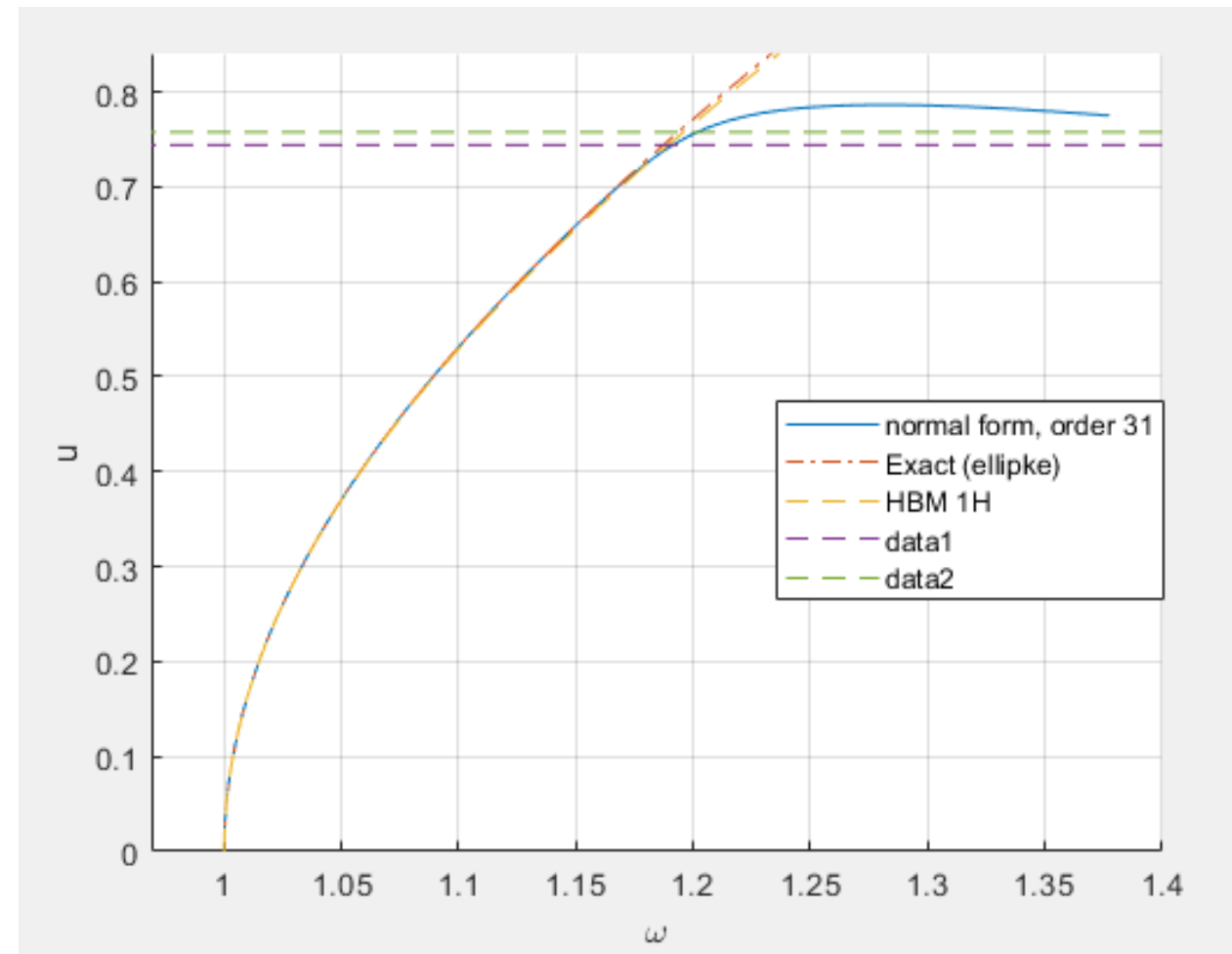
- D'Alembert Criteria

$$\frac{1}{\rho^*} = \lim \frac{a_{n+1}}{a_n}$$

- Cauchy Criteria

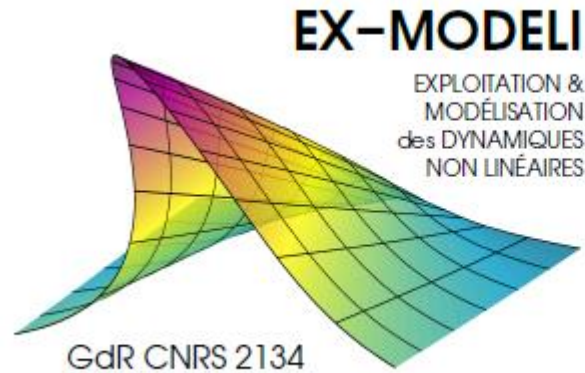
$$\frac{1}{\rho^*} = \lim (a_n)^{1/n}$$

## Duffing example



# Conclusion

- Normal form on quadratic DAE
  - Very general framework for non linear dynamics
  - Use of vector space structure to write the equation explicitly
  - Allows for reduced order model construction
  - Application on geometrically exact FEM (with unit quaternions)
  - Convergence radius (a posteriori validity bound, using Cauchy or D'Alembert criteria)
- Future work
  - Start from the hamiltonian directly (Canonical transformation, Birkoff normal form)
  - Link with the asymptotic numerical method (MANLAB)
  - Link with the Koopman theory



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# Polynomials and multiplication

- Consider for example polynomial of degree  $d = 2$  in  $n = 2$  variables  $(z_1, z_2)$  :

$$p(z) = p_0 + (p_1 z_1 + p_2 z_2) + (p_3 z_1^2 + p_4 z_1 z_2 + p_5 z_2^2) = p_0 + p_1 b_1 + p_2 b_2 + p_3 b_3 + p_4 b_4 + p_5 b_5$$

- Vector space of dimension  $M = 6$
- Basis :  $B = B(z) = (1 \quad z_1 \quad z_2 \quad z_1^2 \quad z_1 z_2 \quad z_2^2)$
- Multiplication:

$$p(z)q(z) = \sum_m p_m b_m \sum_s q_s b_s = \sum_{ms} p_m q_s b_m b_s$$

$$p(z)q(z) = \sum_r \left( \sum_{ms} p_m q_s \Lambda_{ms}^r \right) b_r$$

$$p(z)q(z) = \sum_r (pq)_r b_r$$

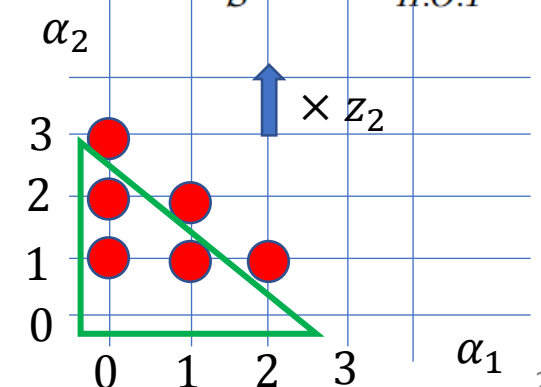
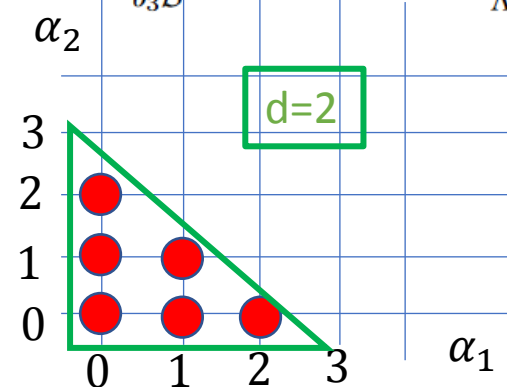
- One only needs to compute the products  $b_m b_s$  and express it relative to the basis:

$$b_m b_s = \Lambda_{ms}^r b_r$$

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- Example :  $b_3 B^T = \Lambda_3 B^T$

$$\underbrace{\begin{pmatrix} z_2 \\ z_1 z_2 \\ z_2^2 \\ z_1^2 z_2 \\ z_1 z_2^2 \\ z_2^3 \end{pmatrix}}_{b_3 B} = \underbrace{\begin{pmatrix} 0 & 0 & 1 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 0 & 0 & 1 \\ 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 \end{pmatrix}}_{\Lambda^3} \underbrace{\begin{pmatrix} 1 \\ z_1 \\ z_2 \\ z_1^2 \\ z_1 z_2 \\ z_2^2 \end{pmatrix}}_B + \underbrace{\begin{pmatrix} 0 \\ 0 \\ 0 \\ z_1^2 z_2 \\ z_1 z_2^2 \\ z_2^3 \end{pmatrix}}_{H.O.T}$$





# Polynomials and derivatives

- Polynomial of degree  $d = 2$  in  $n = 2$  variables  $(z_1, z_2)$  :      • Example :  $\frac{\partial B^T}{\partial z_2} = \nabla_2 B^T$  :

$$p(z) = \sum p_m b_m$$

- Basis :  $B = B(z) = (1 \quad z_1 \quad z_2 \quad z_1^2 \quad z_1 z_2 \quad z_2^2)$

- Derivative:

$$\frac{\partial p}{\partial z_k}(z) = \sum_m p_m \frac{\partial b_m}{\partial z_k} = \sum_m p_m \sum_r \nabla_{mk}^r b_r$$

$$\frac{\partial p}{\partial z_k}(z) = \sum_r \left( \sum_m p_m \nabla_{mk}^r \right) b_r$$

$$\frac{\partial p}{\partial z_k}(z) = \sum_r \left( \frac{\partial p}{\partial z_k} \right)_r b_r$$

- One only needs to compute the derivative of the basis monomials

$$\frac{\partial b_m}{\partial z_k} = \nabla_{mk}^r b_r$$

$$\underbrace{\begin{pmatrix} 0 \\ 0 \\ 1 \\ 0 \\ z_1 \\ 2z_2 \end{pmatrix}}_{\frac{\partial B}{\partial z_2}} = \underbrace{\begin{pmatrix} 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 \\ 1 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 & 0 & 0 \\ 0 & 0 & 2 & 0 & 0 & 0 \end{pmatrix}}_{\nabla^2} \underbrace{\begin{pmatrix} 1 \\ z_1 \\ z_2 \\ z_1^2 \\ z_1 z_2 \\ z_2^2 \end{pmatrix}}_B$$