REDUCTION METHODS IN NONLINEAR

DYNAMICS:

CENTER MANIFOLDS AND NORMAL FORMS

Mariana Haragus

Institut FEMTO-ST, Besançon, France

Journées annuelles du GDR EX-MODELI

Besançon, 9-10 November 2023





MOTIVATION

study nonlinear waves and patterns

NONLINEAR WAVES ON WATER



Water wave [David Sanger Photography]



Solitary wave Lagoon of Molokai, Hawaii [photo: R.I. Odom]



Roll wave Channel in Lions Bay, Canada [website of N. Balmforth]



Mascaret de St Pardon Dordogne river



Tsunami in Asia



Rogue wave Chemical tanker ship Stolt Surf [photo: K. Petersen]

OTHER NONLINEAR WAVES



Kelvin-Helmholtz clouds Mount Duval, Australia [English Wikipedia: GRAHAMUK]



Morning Glory cloud near Burketown, Australia [author: Mick Petrov]



Hurricane



Fire rainbow Northern Idaho



Sound wave Bell Telephone Laboratories [book by David C. Knight]

PATTERNS IN NATURE



Sand patterns [photo: R. Niebrugge]













The Mathematics of ... Nonlinear Waves and Patterns

observed in nature, experiments, numerical simulations

- particular solutions of PDEs or ODEs
 - well-defined temporal and spatial structure
 - e.g., traveling waves

play a key role in the dynamics of the underlying system

The Mathematics of ... Nonlinear Waves and Patterns

Questions

- existence spatial and temporal properties
- **stability** *spatial and temporal behavior*
- interactions
- **•** ...
- role in the dynamics of the system

The Mathematics of ... Nonlinear Waves and Patterns

Answers

- ... many different methods ...
- ... not enough ...

- numerical
- analytical

TWO REDUCTION METHODS

- center manifolds
- normal forms

Some applications ¹

bifurcations of nonlinear waves and patterns

¹Focus on results not on equations!

WATER WAVES



WATER WAVES



WATER WAVES



WATER-WAVE PROBLEM



gravity-capillary water waves

- three-dimensional inviscid fluid layer
- constant density ρ
- gravity and surface tension
- (ir)rotational flow



[Nekrasov, Levi-Civita, Struik, Lavrentiev, Friedrichs & Hyers, ... Amick, Kirchgässner, Iooss, Buffoni, Groves, Toland, Lombardi, Sun, ...]

3D TRAVELING WAVES



[Groves, Mielke, Craig, Nicholls, H., Kirchgässner, Deng, Sun, Sandstede, looss, Plotnikov, Wahlén, ...]

Defects in striped patterns



- grain boundaries
- dislocations
- disclinations

[D. Boyer, J. Viñals]

Defects in Striped Patterns

• Occur in a wide range of systems

- Rayleigh-Bénard convection experiment
- crystal patterns in material science
- chemical reactions
- biology
-

EXISTENCE OF DEFECTS

grain boundaries and dislocations

(Rayleigh-Bénard convection, Swift-Hohenberg equation)





[H., Scheel, Wu, looss, Buffoni, Lloyd, ...]

More defects

□ Some may be treatable by related methods ...







More defects

Some may be treatable by related methods ...







Some cannot be treated by any of these methods ...





A PROBLEM FROM OPTICS

Home-made whispering gallery modes resonators





[Yanne Chembo, Rémi Henriet,

Aurelien Coillet]



FREQUENCY COMBS

 optical signals: superposition of modes with equally spaced frequencies + stationary in suitable reference frame



 analytical results are in very good agreement with numerical and experimental results

[Chembo, Godey, H., Delcey, Reichel, Mandel, ...]

reduce dimensions

Dynamical system (*infinite-dimensional*)

$$rac{\mathsf{d}}{\mathsf{d}\mathsf{t}} \, \mathsf{U} = \mathsf{F}(\mathsf{U}), \quad \mathsf{U}(t) \in \mathcal{X}$$

■ start with particular solution: equilibrium

$$\mathbf{U}(t) = \mathbf{U}_{*}$$

(often $\mathbf{U}_* = 0$, but not always)

Question: local dynamics?

EXAMPLE

Swift-Hohenberg equation:

$$\frac{\partial u}{\partial t} = -\left(1 + \frac{\partial^2}{\partial x^2}\right)^2 u + \mu u - u^3$$

• u(x, t) 2 π -periodic in x, parameter $\mu \in \mathbb{R}$

EXAMPLE

Swift-Hohenberg equation:

$$\frac{\partial u}{\partial t} = -\left(1 + \frac{\partial^2}{\partial x^2}\right)^2 u + \mu u - u^3$$

• u(x,t) 2 π -periodic in x, parameter $\mu \in \mathbb{R}$

Dynamical system:

$$rac{\mathsf{d}}{\mathsf{d}\mathsf{t}}\,\mathsf{U}=\mathsf{F}\,(\mathsf{U},\mu),\quad \mathit{U}(t)\in\mathcal{X}$$

•
$$\mathbf{U} = u$$
, $\mathbf{F}(\mathbf{U}, \mu) = -\left(1 + \frac{\partial^2}{\partial x^2}\right)^2 u + \mu u - u^3$

•
$$\mathcal{X} = L^2_{per}(0, 2\pi)$$

• particular solution: $\mathbf{U}_* = \mathbf{0}$

LINEAR DYNAMICS

Linearized equation

$$\frac{\mathrm{d}}{\mathrm{d} t} \mathsf{U} = \mathcal{A}_* \mathsf{U}, \qquad \mathcal{A}_* \mathsf{U} = D_{\mathsf{U}} \mathcal{F}(\mathsf{U}_*)$$

• spectrum of A_* :



LINEAR DYNAMICS

Linearized equation

$$\frac{\mathrm{d}}{\mathrm{d} t} \mathrm{U} = \mathcal{A}_* \mathrm{U}, \qquad \mathcal{A}_* \mathrm{U} = D_{\mathrm{U}} \mathcal{F}(\mathrm{U}_*)$$





- center space X_c: sum of the generalized eigenspaces associated with purely imaginary eigenvalues
- X_c contains all bounded solutions

NONLINEAR DYNAMICS



$$\frac{d}{dt} U = \mathcal{A}_* U + \mathcal{R}(U)$$

Center manifold

- \blacksquare analogue of the center space \boldsymbol{X}_{c} for the nonlinear equation
- contains all small bounded solutions of the dynamical system
- exist in finite dimensions
- infinite dimensions: three main hypotheses

Hypothesis

1 \mathcal{A}_* is a closed operator in a Hilbert (Banach) space \mathcal{X} with dense domain $\mathcal{Y} \subset \mathcal{X}$; $\mathcal{R} : \mathcal{Y} \to \mathcal{Z}$ is well defined.

Hypothesis

- A_{*} is a closed operator in a Hilbert (Banach) space X with dense domain Y ⊂ X; R : Y → Z is well defined.
- spectrum of A_{*}: finite number of purely imaginary eigenvalues



Hypothesis

- A_{*} is a closed operator in a Hilbert (Banach) space X with dense domain Y ⊂ X; R : Y → Z is well defined.
- spectrum of A_{*}: finite number of purely imaginary eigenvalues
- **3** resolvent estimates:

(see also [H. & looss, 2011])

$$\|(\mathcal{A}_* - i\omega)^{-1}\|_{\mathcal{X} \to \mathcal{X}} \leq \frac{C}{|\omega|}, \quad \|(\mathcal{A}_* - i\omega)^{-1}\|_{\mathcal{Z} \to \mathcal{Y}} \leq \frac{C}{|\omega|^{1-\alpha}}$$

for $|\omega| \ge \omega_*$, and some $\alpha \in [0, 1)$.

Theorem

The dynamical system

$$\frac{\mathrm{d}\mathsf{U}}{\mathrm{d}t} = \mathcal{A}_* + \mathcal{R}(\mathsf{U})$$
 possesses a locally

invariant manifold

 $\mathcal{M}_c = \{\mathsf{U} = \mathsf{U}_0 + \Phi(\mathsf{U}_0) ; \ \mathsf{U}_0 \in \mathsf{X}_c\}$

Theorem

The dynamical system

invariant manifold

$$\frac{\mathrm{dU}}{\mathrm{dt}} = \mathcal{A}_* + \mathcal{R}(\mathrm{U})$$
 possesses a locally

$$\mathcal{M}_c = \{\mathsf{U} = \mathsf{U}_0 + \Phi(\mathsf{U}_0) ; \ \mathsf{U}_0 \in X_c\}$$

• X_c is the (finite-dimensional) center space of the linearized equation $X_c = \bigoplus_{i \in I} E_{i \in I}$

$$X_c = \bigoplus_{i\kappa\in\sigma(\mathcal{A}_*)} E_{i\kappa};$$

• ϕ is a map of class C^k ;

• \mathcal{M}_c contains all bounded solutions of the system.

Reduced dynamics:

 solutions of the infinite-dimensional dynamical system which belong to the center manifold

 $\mathsf{U}(t) = \mathsf{U}_{\mathbf{0}}(t) + \mathbf{\Phi}(\mathsf{U}_{\mathbf{0}}(t)), \quad \mathsf{U}_{\mathbf{0}}(t) \in \mathsf{X}_{\mathsf{c}}$

■ **U**₀(*t*) solves the reduced system

$$\frac{\mathsf{d} \mathsf{U}_{\mathbf{0}}}{\mathsf{d} t} = \mathcal{A}_{0} \mathsf{U}_{0} + \mathcal{R}_{0}(\mathsf{U}_{\mathbf{0}})$$

(the dimension of this system is often small)

• $\mathcal{A}_0 = \mathcal{A}_*|_{\mathbf{X}_c}$ and the Taylor expansion of $\mathcal{R}_0(U_0)$ can be computed ...

- locally invariant manifolds tangent to X_c at 0;
 - their dimension is equal to the dimension of X_c
- not unique, but the Taylor expansion of **Φ** is unique
- manifolds of class C^k for C^k vector fields, but not analytic ...

FURTHER PROPERTIES

- parameter-dependent version
- symmetries of the original system are inherited by the reduced system, e.g.,
 - equivariance
 - reversibility
 - Hamiltonian structure

- open questions . . .
 - *infinite-dimensional center manifolds*?
 - quasilinear dynamical systems in Banach spaces?
- used in the study of local bifurcations (existence of nonlinear waves) / stability problems
- together with stable/unstable manifolds provide a description of the nonlinear dynamics
- together with normal forms provide a rigorous justification of amplitude equations

NORMAL FORMS

simplify nonlinear vector fields

NORMAL FORMS

Dynamical system (finite-dimensional)

$$\frac{du}{dt} = \mathbf{L}u + \mathbf{R}(u) \tag{1}$$

Hypotheses

1
$$L \in \mathcal{L}(\mathbb{R}^n);$$

2 for $k \ge 2$, there exists a neighborhood $\mathcal{V} \subset \mathbb{R}^n$ of 0 such that $\mathbf{R} \in \mathcal{C}^k(\mathcal{V}, \mathbb{R}^n)$ and

$$\mathbf{R}(0)=0,\quad D\mathbf{R}(0)=0.$$

THEOREM

There exists a change of variables $u = v + \Phi(v)$, with $\Phi : \mathbb{R}^n \to \mathbb{R}^n$ a polynomial of degree *p*, which transforms the system (1) in its "normal form"

$$\frac{dv}{dt} = \mathbf{L}v + \mathbf{N}(v) + \boldsymbol{\rho}(v)$$

with the following properties:

THEOREM

There exists a change of variables $u = v + \Phi(v)$, with $\Phi : \mathbb{R}^n \to \mathbb{R}^n$ a polynomial of degree *p*, which transforms the system (1) in its "normal form"

$$\frac{dv}{dt} = \mathbf{L}v + \mathbf{N}(v) + \boldsymbol{\rho}(v)$$

with the following properties:

- ρ is of class C^k in a neighborhood of 0, and $\rho(v) = o(||v||^p)$;
- $N : \mathbb{R}^n \to \mathbb{R}^n$ is a polynomial of degre p;

Properties of $\boldsymbol{\mathsf{N}}$

1
$$\mathbf{N}: \mathbb{R}^n \to \mathbb{R}^n$$
 polynomial of degre p , $|\mathbf{N}(0) = 0, D\mathbf{N}(0) = 0$

2
$$\mathsf{N}(e^{t\mathsf{L}^*}v) = e^{t\mathsf{L}^*}\mathsf{N}(v)$$
 for all $(t, v) \in \mathbb{R} \times \mathbb{R}^n$

3
$$DN(v)L^*v = L^*N(v)$$
 for all $v \in \mathbb{R}^n$

we use the identity

$$\frac{d}{dt}\left(e^{-t\mathsf{L}^*}\mathsf{N}(e^{t\mathsf{L}^*}v)\right) = e^{-t\mathsf{L}^*}\left(-\mathsf{L}^*\mathsf{N}(e^{t\mathsf{L}^*}v) + D\mathsf{N}(e^{t\mathsf{L}^*}v)\mathsf{L}^*e^{t\mathsf{L}^*}v\right)$$

[C. Elphick et al.]

•
$$\frac{du}{dt} = \mathbf{L}u + \mathbf{R}(u), \quad u = v + \mathbf{\Phi}(v), \quad \frac{dv}{dt} = \mathbf{L}v + \mathbf{N}(v) + \rho(v)$$

 $(\mathbb{I} + D\mathbf{\Phi}(v)) (\mathbf{L}v + \mathbf{N}(v) + \rho(v)) = \mathbf{L}(v + \mathbf{\Phi}(v)) + \mathbf{R}(v + \mathbf{\Phi}(v))$

•
$$\frac{du}{dt} = \mathbf{L}u + \mathbf{R}(u), \quad u = v + \mathbf{\Phi}(v), \quad \frac{dv}{dt} = \mathbf{L}v + \mathbf{N}(v) + \rho(v)$$

 $(\mathbb{I} + D\mathbf{\Phi}(v)) (\mathbf{L}v + \mathbf{N}(v) + \boldsymbol{\rho}(v)) = \mathbf{L}(v + \boldsymbol{\Phi}(v)) + \mathbf{R}(v + \boldsymbol{\Phi}(v))$

•
$$\mathbf{R}(u) = \sum_{2 \le q \le p} \mathbf{R}_q(u^{(q)}) + o(||u||^p),$$

• $\mathbf{\Phi}(v) = \sum_{2 \le q \le p} \mathbf{\Phi}_q(v^{(q)}), \ \mathbf{N}(v) = \sum_{2 \le q \le p} \mathbf{N}_q(v^{(q)})$

$$D\Phi_2(v^{(2)})\mathsf{L}v - \mathsf{L}\Phi_2(v^{(2)}) = \mathsf{R}_2(v^{(2)}) - \mathsf{N}_2(v^{(2)})$$

$$D\mathbf{\Phi}_q(\mathbf{v}^{(q)})\mathbf{L}\mathbf{v} - \mathbf{L}\mathbf{\Phi}_q(\mathbf{v}^{(q)}) = \mathbf{Q}_q(\mathbf{v}^{(q)}) - \mathbf{N}_q(\mathbf{v}^{(q)})$$

•
$$\frac{du}{dt} = \mathbf{L}u + \mathbf{R}(u), \quad u = v + \mathbf{\Phi}(v), \quad \frac{dv}{dt} = \mathbf{L}v + \mathbf{N}(v) + \rho(v)$$

 $(\mathbb{I} + D\mathbf{\Phi}(v)) (\mathbf{L}v + \mathbf{N}(v) + \boldsymbol{\rho}(v)) = \mathbf{L}(v + \boldsymbol{\Phi}(v)) + \mathbf{R}(v + \boldsymbol{\Phi}(v))$

•
$$\mathbf{R}(u) = \sum_{2 \le q \le p} \mathbf{R}_q(u^{(q)}) + o(||u||^p),$$

 $\mathbf{\Phi}(v) = \sum_{2 \le q \le p} \mathbf{\Phi}_q(v^{(q)}), \ \mathbf{N}(v) = \sum_{2 \le q \le p} \mathbf{N}_q(v^{(q)})$

$$D\Phi_2(v^{(2)})\mathsf{L}v - \mathsf{L}\Phi_2(v^{(2)}) = \mathsf{R}_2(v^{(2)}) - \mathsf{N}_2(v^{(2)})$$

$$D\mathbf{\Phi}_q(\mathbf{v}^{(q)})\mathbf{L}\mathbf{v} - \mathbf{L}\mathbf{\Phi}_q(\mathbf{v}^{(q)}) = \mathbf{Q}_q(\mathbf{v}^{(q)}) - \mathbf{N}_q(\mathbf{v}^{(q)})$$

• Linear equation: $A_L \Phi_q = Q_q - N_q$

• Solve the linear equation:

$$\mathcal{A}_{\mathsf{L}} \mathbf{\Phi}_q = \mathbf{Q}_q - \mathbf{N}_q$$

- $\bullet \ \mathcal{A}_{\mathsf{L}^*} \text{ adjoint of } \mathcal{A}_{\mathsf{L}}$
- $\mathbf{Q}_q \mathbf{N}_q \in \ker(\mathcal{A}_{\mathbf{L}^*})^{\perp} = \operatorname{im}(\mathcal{A}_{\mathbf{L}})$

$$\bullet \mathbf{P}_{\ker(\mathcal{A}_{\mathsf{L}^*})}(\mathbf{Q}_q - \mathbf{N}_q) = 0$$

$$\blacksquare \ \mathsf{N}_q = \mathsf{P}_{\ker(\mathcal{A}_{\mathsf{L}^*})} \mathsf{Q}_q \in \ker(\mathcal{A}_{\mathsf{L}^*})$$

Remark: the normal form is not unique

EXAMPLES

• 0² normal form

(Takens-Bogdanov bifurcation)

$$\mathbf{L} = \begin{pmatrix} 0 & 1 \\ 0 & 0 \end{pmatrix}, \quad \mathbf{v} = \begin{pmatrix} A \\ B \end{pmatrix} \in \mathbb{R}^2, \qquad \mathbf{N}(u) = \begin{pmatrix} AP(A) \\ BP(A) + Q(A) \end{pmatrix}$$
$$P(0) = Q(0) = Q'(0) = 0$$

EXAMPLES

■ 0² normal form

(Takens-Bogdanov bifurcation)

$$\mathbf{L} = \begin{pmatrix} 0 & 1 \\ 0 & 0 \end{pmatrix}, \quad \mathbf{v} = \begin{pmatrix} A \\ B \end{pmatrix} \in \mathbb{R}^2, \quad \mathbf{N}(u) = \begin{pmatrix} AP(A) \\ BP(A) + Q(A) \end{pmatrix}$$

P(0) = Q(0) = Q'(0) = 0

• $i\omega$ normal form

(Hopf bifurcation)

$$\mathbf{L} = \begin{pmatrix} i\omega & 0\\ 0 & -i\omega \end{pmatrix}, \quad \mathbf{v} = \begin{pmatrix} A\\ \overline{A} \end{pmatrix}, \quad \mathbf{N}(\mathbf{v}) = \begin{pmatrix} AQ(|A|^2)\\ \overline{AQ}(|A|^2) \end{pmatrix}$$

 $Q:\mathbb{C}\to\mathbb{C},\quad Q(0)=0$

FURTHER PROPERTIES

- parameter-dependent version
- symmetries of the original system are inherited by the reduced system, e.g.,
 - equivariance
 - reversibility

Nonlinear waves and patterns



Approach

spatial dynamics

IDEA OF SPATIAL DYNAMICS



Klaus Kirchgässner (1931 – 2011)

SPATIAL DYNAMICS

nonperiodic bounded solutions of PDEs in infinite strips



[Kirchgässner, 1982]

• x timelike coordinate

SPATIAL DYNAMICS

nonperiodic bounded solutions of PDEs in infinite strips



[Kirchgässner, 1982]

• x timelike coordinate

Dynamical system

$$rac{\mathsf{d}}{\mathsf{d}\mathsf{x}}\mathsf{U}=\mathsf{F}(\mathsf{U},\mu),\quad \mathsf{U}(\mathsf{x})\in\mathcal{X}$$

- U(x) belongs to a Hilbert (Banach) space X of functions depending upon the "space" variables;
- $\mu \in \mathbb{R}^m$ parameters.

EXAMPLE

$$\textbf{u}_{\textbf{x}\textbf{x}} + \textbf{u}_{\textbf{y}\textbf{y}} = \textbf{f}(\textbf{u}), \quad (\textbf{x},\textbf{y}) \in \mathbb{R} \times (\textbf{0},\textbf{1})$$

• Set
$$v = u_x$$

Dynamical system:

$$\frac{\mathrm{d}}{\mathrm{dx}} \mathrm{U} = \mathrm{F}(\mathrm{U}), \quad \mathrm{U} = \begin{pmatrix} u \\ v \end{pmatrix}, \quad \mathrm{F}(\mathrm{U}) = \begin{pmatrix} v \\ -u_{yy} + f(u) \end{pmatrix}$$

• Phase space: $\boldsymbol{U}(\boldsymbol{x}) \in \boldsymbol{\mathcal{X}}, \ \boldsymbol{\mathcal{X}} = L^2(0,1) \times L^2(0,1), \dots$

SPATIAL DYNAMICS APPROACH

1 Dynamical system

$$rac{\mathsf{d}}{\mathsf{d}\mathsf{x}} \, \mathsf{U} = \mathsf{F} \, (\mathsf{U}, \mu), \quad \mathsf{U}(\mathsf{x}) \in \mathcal{X}$$

look for bounded solutions

Spatial Dynamics Approach

1 Dynamical system

$$rac{\mathsf{d}}{\mathsf{d}\mathsf{x}}\,\mathsf{U}=\mathsf{F}\,(\mathsf{U},\mu),\quad \mathsf{U}(\mathsf{x})\in\mathcal{X}$$

look for bounded solutions







REDUCTION

1 Dynamical system

$$rac{\mathsf{d}}{\mathsf{d}\mathsf{x}} \mathsf{U} = \mathsf{F}(\mathsf{U},\mu), \quad \mathsf{U}(\mathsf{x}) \in \mathcal{X}$$

REDUCTION

1 Dynamical system

$$rac{\mathsf{d}}{\mathsf{d}\mathsf{x}}\mathsf{U}=\mathsf{F}(\mathsf{U},\mu),\quad \mathsf{U}(\mathsf{x})\in\mathcal{X}$$

2 Center manifold reduction: obtain a reduced system of

ODEs ...
$$\frac{\mathsf{d}}{\mathsf{d}\mathsf{x}}\,\mathsf{v}=\mathsf{g}\,(\mathsf{v},\mu),\quad \mathsf{v}(\mathsf{x})\in\mathbb{R}^d$$

(pass from infinite to finite dimensions)

REDUCTION

1 Dynamical system

$$rac{\mathsf{d}}{\mathsf{d}\mathsf{x}}\mathsf{U}=\mathsf{F}(\mathsf{U},\mu),\quad \mathsf{U}(\mathsf{x})\in\mathcal{X}$$

2 Center manifold reduction: obtain a reduced system of

ODEs ...
$$\frac{\mathsf{d}}{\mathsf{d}\mathsf{x}}\,\mathsf{v}=\mathsf{g}\,(\mathsf{v},\mu),\quad \mathsf{v}(\mathsf{x})\in\mathbb{R}^d$$

(pass from infinite to finite dimensions)





REDUCED SYSTEM

- 1 Dynamical system
- **2** Center manifold reduction: reduced system of ODEs

$$rac{\mathsf{d}}{\mathsf{d}\mathsf{x}}\mathsf{v}=\mathsf{g}\,(\mathsf{v},\mu),\quad \mathsf{v}(\mathsf{x})\in\mathbb{R}^d$$

- **③** Bounded orbits of the reduced system of ODEs
 - e.g., use normal forms
 - study a truncated system, then show persistence of the truncated dynamics

REDUCED SYSTEM

- 1 Dynamical system
- **2** Center manifold reduction: reduced system of ODEs

$$rac{\mathsf{d}}{\mathsf{d}\mathsf{x}}\mathsf{v}=\mathsf{g}\,(\mathsf{v},\mu),\quad \mathsf{v}(\mathsf{x})\in\mathbb{R}^d$$

- **③** Bounded orbits of the reduced system of ODEs
 - e.g., use normal forms
 - study a truncated system, then show persistence of the truncated dynamics
 - difficult:





Nonlinear waves and patterns



Computation of the reduced

SYSTEM

Taylor expansion of

$$\mathcal{R}_0(\mathsf{U}_0) = \mathsf{P}_c \mathcal{R}(\mathsf{U}_0 + \Phi(\mathsf{U}_0))$$

1 \mathcal{R} known

2 $U_0(x)$ belongs to the center space $X_c = \bigoplus_{i \kappa \in \sigma(A_*)} E_{i\kappa}$

- U₀(x) is a finite linear combination of basis vectors
- basis of X_c: consists of generalized eigenvectors which form a basis for the generalized eigenspaces E_{iκ}
- **3** P_c is the spectral projector onto tel que
 - $P_c^2 = P_c$ (projector)
 - $P_c A_* = A_* P_c$ (commutes with A_*)
- Φ : it is possible to compute its Taylor expansion (often, not necessary)