

ÉCOLE DOCTORALE Sciences mécaniques et énergétiques, matériaux et géosciences (SMEMAG)

Direct parametrisation of invariant manifold with shell finite element: nonlinear dynamics of thin structures using reduced order modeling

Abundant application scenarios of thin structures

Where do we stand?

Normal form method

- Touzé C. A **normal form approach** for nonlinear normal modes
- Vizzaccaro A, et al. **Direct computation** of nonlinear mapping via **normal form** for reduced-order models of finite element nonlinear structures

DPIM (Direct parametrisation of invariant manifold)

- Opreni A, et al. High-order **direct parametrisation of invariant manifolds** for model order reduction of finite element structures: application to generic forcing terms and parametrically excited systems
- Vizzaccaro A, et al. Direct parametrisation of invariant manifolds for generic **non-autonomous** systems including superharmonic resonances

batteries

Shell finite element modeling

- **Fewer degrees of freedom** compared to solid elements
- More convenient definition of the **kinematic description** of thin structures in the transverse direction
- Introduce **assumed natural strain** to prevent Poisson locking

Nonlinear dynamics solution methods

- Increment harmonic balance method (For **full order models**)
- Collocation method (For **reduced order models**)

These theories lay the foundation for our development of a framework for solving thin structures !

Curved shell structure modeling: governing equations

Positional relationship

$$
\mathbf{X}(\theta^{\alpha}, \theta^3) = \mathbf{R}(\theta^{\alpha}) + \theta^3 \mathbf{a}_3(\theta^{\alpha}), \alpha = 1, 2
$$

$$
\mathbf{x}(\theta^{\alpha}, \theta^3) = \mathbf{r}(\theta^{\alpha}) + \theta^3 \tilde{\mathbf{a}}_3(\theta^{\alpha}), \alpha = 1, 2
$$

Covariant base tensor

$$
\begin{aligned} &\mathbf{G}_{\alpha} = \mathbf{X}_{\alpha} = \mathbf{a}_{\alpha} + \theta^3 \mathbf{a}_{3,\alpha} \quad \mathbf{g}_{\alpha} = \mathbf{x}_{,\alpha} = \tilde{\mathbf{a}}_{\alpha} + \theta^3 \tilde{\mathbf{a}}_{3,\alpha} \\ &\mathbf{G}_{3} = \mathbf{X}_{3} = \mathbf{a}_{3} \qquad \qquad \mathbf{g}_{3} = \mathbf{x}, 3 = \tilde{\mathbf{a}}_{3} \end{aligned}
$$

Green-Lagrange strain

$$
E_{ij} = \frac{1}{2} \left(\mathbf{g}_i \cdot \mathbf{g}_j - \mathbf{G}_i \cdot \mathbf{G}_j \right)
$$

$$
E_{ij} = \frac{1}{2} \left(\mathbf{G}_i \cdot \frac{\partial \mathbf{u}}{\partial \theta^j} + \mathbf{G}_j \cdot \frac{\partial \mathbf{u}}{\partial \theta^i} + \frac{\partial \mathbf{u}}{\partial \theta^i} \frac{\partial \mathbf{u}}{\partial \theta^j} \right)
$$

Constitutive relation

Traditional shell elements lead to locking issues $E_{33}^{(0)} + \theta^3 E_{33}^{(1)} \simeq -\; \frac{D^{33ij}}{D^{3333}}(E_{ij}^{(0)} + \theta^3 E_{ij}^{(1)}) \; .$

Extend traditional variational equations to avoid locking issues

Curved shell structure modeling: finite element discretisation

Shape function interpolation

Green-Lagrange strain

$$
\left\{ \mathbf{E} \right\} = \left(\left[\mathbf{B}_l \right] + \frac{1}{2} \left[\mathbf{B}_{nl}(\mathbf{u}) \right] \right) \left\{ \mathbf{u}^{(e)} \right\} = \left(\left[\mathbf{R} \right] + \frac{1}{2} \left[\mathbf{A}(\mathbf{u}) \right] \right) \left[\mathbf{\Xi} \right] \left\{ \mathbf{u}^{(e)} \right\}
$$

 $\{\delta \mathbf{E}\} = \left(\left[\mathbf{B}_{l}\right] + \left[\mathbf{B}_{nl}(\mathbf{u})\right]\right)\left\{\mathbf{u}^{(e)}\right\} = \left(\left[\mathbf{R}\right] + \left[\mathbf{A}(\mathbf{u})\right]\right)\left[\mathbf{\Xi}\right]\left\{\delta \mathbf{u}^{(e)}\right\}$

Interpolation of enhanced assumed strain

$$
\tilde{E} = [\mathbf{B}_{\alpha}] \{ \alpha^{(e)} \} \text{ with } \tilde{E}_{33} = \alpha_{1} + \alpha_{2} \theta^{1} + \alpha_{3} \theta^{2} + \alpha_{4} \theta^{1} \theta^{2}
$$
\n
$$
N^{i} = \frac{1}{4} (1 + \theta^{1} \theta_{i}^{1}) (1 + \theta^{2} \theta_{i}^{2}) (\theta^{1} \theta_{i}^{1} + \theta^{2} \theta_{i}^{2} - 1) \text{ with } i = 1, 2, 3, 4
$$
\n
$$
N^{i} = \frac{1}{2} (1 - (\theta^{1})^{2}) (1 + \theta^{2} \theta_{i}^{2}) \text{ with } i = 5, 7
$$
\n
$$
N^{i} = \frac{1}{2} (1 + \theta^{1} \theta_{i}^{1}) (1 - (\theta^{2})^{2}) \text{ with } i = 6, 8
$$
\n
$$
\theta^{2} = -1
$$
\n
$$
\theta^{3} = +1
$$
\n
$$
\theta^{4} = +1
$$
\n
$$
\theta^{5} = +1
$$
\n
$$
\theta^{6} = +1
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\theta^{7} = +1
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\theta^{8} = +1
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$$
\theta^{9} = +1
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\n
$$
\theta^{1} = +1
$$
\n
$$
\theta^{2} = +1
$$
\n
$$
\theta^{3} = -1
$$
\n
$$
\theta^{2} = +1
$$

Development of the dynamic equation

 $\{{\bf u}\}=[{\bf N}]\{{\bf u}^{(e)}\}, \quad \{\delta {\bf u}\}=[{\bf N}]\{\delta {\bf u}^{(e)}\}, \quad \{\nabla {\bf u}\}=[\boldsymbol{\Xi}]\{{\bf u}^{(e)}\}\Big|\, [{\bf M}]\{\ddot{{\bf U}}\} + [{\bf C}]\{\dot{{\bf U}}\} + [{\bf K}_l\{{\bf U}\} + \{{\bf G}({\bf U},{\bf U})\} + \{{\bf H}({\bf U},{\bf U},{\bf U})\} = \{{\bf F}(t)\}\Big|$

where

$$
\begin{split}\n\left[\mathbf{M}\right] &= \bigwedge_{k=1}^{N} \int_{\Omega} \rho\left[\mathbf{N}\right]^{T} \left[\mathbf{N}\right] d\Omega \\
\left[\mathbf{K}_{l}\right] &= \bigwedge_{k=1}^{N} \left(\int_{\Omega} \left[\mathbf{B}_{l}\right]^{T} \left[\mathbf{D}\right] \left[\mathbf{B}_{l}\right] d\Omega - \int_{\Omega} \left[\mathbf{B}_{l}\right]^{T} \left[\mathbf{D}\right] \left[\mathbf{B}_{\alpha}\right] d\Omega \times \left[\mathbf{k}_{\alpha\alpha}\right]^{-1} \times \int_{\Omega} \left[\mathbf{B}_{\alpha}\right]^{T} \left[\mathbf{D}\right] \left[\mathbf{B}_{l}\right] d\Omega\right) \\
\left\{ \mathbf{G}(\mathbf{U},\mathbf{U}) \right\} &= \bigwedge_{k=1}^{N} \left(\int_{\Omega} \left[\mathbf{B}_{nl}(\mathbf{u})\right]^{T} \left[\mathbf{D}\right] \left[\mathbf{B}_{l}\right] d\Omega - \frac{1}{2} \int_{\Omega} \left[\mathbf{B}_{l}\right]^{T} \left[\mathbf{D}\right] \left[\mathbf{B}_{nl}(\mathbf{u})\right] d\Omega \\
&- \frac{1}{2} \int_{\Omega} \left[\mathbf{B}_{l}\right]^{T} \left[\mathbf{D}\right] \left[\mathbf{B}_{\alpha}\right] d\Omega \times \left[\mathbf{k}_{\alpha\alpha}\right]^{-1} \times \int_{\Omega} \left[\mathbf{B}_{\alpha}\right]^{T} \left[\mathbf{D}\right] \left[\mathbf{B}_{nl}(\mathbf{u})\right] d\Omega \\
&- \int_{\Omega} \left[\mathbf{B}_{nl}(\mathbf{u})\right]^{T} \left[\mathbf{D}\right] \left[\mathbf{B}_{\alpha}\right] d\Omega \times \left[\mathbf{k}_{\alpha\alpha}\right]^{-1} \times \int_{\Omega} \left[\mathbf{B}_{\alpha}\right]^{T} \left[\mathbf{D}\right] \left[\mathbf{B}_{l}\right] d\Omega\right) \left\{ \mathbf{u}^{(e)} \right\} \\
\left\{ \mathbf{H}(\mathbf{U},\mathbf{U},\mathbf{U}) \right\} &= \bigwedge_{k=1}^{N} \left(\frac{1}{2} \int_{
$$

Establish standard nonlinear dynamic equations to address DPIM solutions

Direct parametrisation of invariant manifold

Second-order dynamic equation First-order dynamic equation $\mathbf{B}[\mathbf{y}] = [\mathbf{A}]\{\mathbf{y}\} + \{\mathbf{Q}(\mathbf{y}, \mathbf{y})\} + \{\mathbf{\Upsilon}\}\tilde{z}$ and $\dot{\tilde{z}} = \tilde{\lambda}\tilde{z}$ $[{\bf M}]\{\ddot{{\bf U}}\} + [{\bf C}]\{\dot{{\bf U}}\} + [{\bf K}_i\{{\bf U}\} + \{{\bf G}({\bf U},{\bf U})\} + \{{\bf H}({\bf U},{\bf U},{\bf U})\} = {\bf F}(t)\}$ $\begin{aligned} \mathbf{E}[\mathbf{B}]\,\langle\frac{\partial \mathbf{W}(\mathbf{z})}{\partial \mathbf{z}}\mathbf{f}(\mathbf{z})\,\rangle_{\,p} \!=\! [\mathbf{A}]\langle \mathbf{W}(\mathbf{z})\rangle_{\,p} + \langle \mathbf{Q}\left(\mathbf{W}(\mathbf{z}),\mathbf{W}(\mathbf{z})\right)\rangle_{\,p} + \langle \mathbf{\Upsilon}\tilde{z}\rangle_{\,p} \end{aligned}$ **Nonlinear mapping** $\mathbf{W}(\mathbf{z}) = \sum_{\substack{p=1\o} } \left\langle \mathbf{W}(\mathbf{z}) \right\rangle_p.$ $\{y\} = W(z)$ $\begin{array}{|c|c|c|c|}\hline \textbf{B} & \textbf{B} & \textbf{b} \end{array} \begin{array}{|c|c|c|}\hline \textbf{B} & \textbf{B} & \textbf{b} \end{array} \begin{array}{|c|c|c|}\hline \textbf{B} & \textbf{B} & \textbf{b} \end{array} \begin{array}{|c|c|c|c|}\hline \textbf{B} & \textbf{B} & \textbf{b} \end{array} \begin{array}{|c|c|c|c|}\hline \textbf{B} & \textbf{B} & \textbf{b} \end{array} \begin{array}{|c|c|c|c|$ $\mathbf{f}(\mathbf{z}) = \sum^{o} \left<\mathbf{f}(\mathbf{z})\right>_{p}$ **ROM equations** $\begin{split} \mathbf{E}[\mathbf{B}]\, \langle \frac{\partial \mathbf{W}(\mathbf{z})}{\partial \mathbf{z}}\, \mathbf{f}(\mathbf{z}) \, \rangle_{\,p} \!=\! [\mathbf{A}]\langle \mathbf{W}(\mathbf{z}) \rangle_{\,p} + \langle \boldsymbol{Q}\left(\mathbf{W}(\mathbf{z}), \mathbf{W}(\mathbf{z})\right) \rangle_{\,p}\! \quad p > 1 \, . \end{split}$ $\{\dot{\mathbf{z}}\} = \mathbf{f}(\mathbf{z})$

Order-1 homological equation

$$
\begin{bmatrix} \tilde{\lambda} \mathbf{B} - \mathbf{A} & \mathbf{B} \boldsymbol{\Phi}_{\mathcal{R}} & 0 \\ \boldsymbol{\Phi}_{\mathcal{R}}^{\dagger} \mathbf{B} & 0 & 0 \\ 0 & 0 & \mathbf{I} \end{bmatrix} \begin{bmatrix} \mathbf{W}^{(1,d+1)} \\ \mathbf{f}_{\mathcal{R}}^{(1,d+1)} \\ \mathbf{f}_{\mathcal{K}}^{(1,d+1)} \end{bmatrix} = \begin{bmatrix} \mathbf{\Upsilon} \\ 0 \\ 0 \end{bmatrix}
$$

Order-p homological equation

$$
\begin{bmatrix}\n\sigma^{(p,k)}\mathbf{B}-\mathbf{A} & \mathbf{B}\mathbf{\Phi}_{\mathcal{R}} & 0 \\
\mathbf{\Phi}_{\mathcal{R}}^{\dagger}\mathbf{B} & 0 & 0 \\
0 & 0 & \mathbf{I}\n\end{bmatrix}\n\begin{bmatrix}\n\mathbf{W}^{(p,k)} \\
\mathbf{f}_{\mathcal{R}}^{(p,k)} \\
\mathbf{f}_{\mathcal{R}}^{(p,k)}\n\end{bmatrix} =\n\begin{bmatrix}\n\mathbf{R}^{(p,k)} \\
0 \\
0\n\end{bmatrix}
$$

Numerical example: modal analysis

Numerical example: primary mode

Numerical example: 1:3 superharmonic resonance

 $1.5\,$

1.45

 1.4

1.35

1.25

 $1.2\,$

1.15

(Solid/Shell)

 Φ Scale $1.3\,$

Numerical example: 1:2 superharmonic resonance

Numerical example: 1:2 superharmonic resonance

Highlights

1. Established a **shell element model** based on DPIM for application in reduced-order methods for nonlinear dynamics of thin structures

2. By comparing with solid elements and full-order model computation methods, we verified the **effectiveness of the shell element** developed in this work in reduced-order methods

3. Compared to solid elements, the **computational cost is lower**

4. The **kinematic assumptions** of thin structures can be conveniently applied in the thickness direction

Future works

1. Based on the validation results of this work, we will proceed to study more **practical engineering problems** concerning complex nonlinear geometrically thin structures, which are not convenient to analyze using solid elements

2. On this basis, we will also consider more complex working conditions, including **fluid-structure interaction** problems and **material nonlinearity** issues

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Thanks for your listening

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