# Modelling and Control of Distributed Parameter Systems: A port-Hamiltonian Approach Transfer Functions

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## Introduction

The aim of this part is to define transfer function for systems described by partial differential equations.

We derive these transfer functions via a very simple calculation. For port-Hamiltonian systems we show that the energy/power balance induces properties on the transfer function.

## Transfer function for an o.d.e.

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the transfer function of this system is given by

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- Laplace transform, or
- Exponential solutions.

One way for obtaining the transfer function of

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is to take  $u(t)=e^{st}$ ,  $s\in\mathbb{C}$ , and to try to find a solution of the same format, i.e.,  $y(t)=\alpha e^{st}$ .

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If  $s \neq -5$ , this is solvable;

$$\alpha = \frac{3}{s+5}.$$

So if we want to find an exponential solution

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▶ We call this the transfer function at s.

## Transfer function via exponential solutions

#### Definition

Given an (abstract) differential equation in the variables (u(t),z(t),y(t)), where u(t), z(t), and y(t) take their values in the (Hilbert) spaces U, Z, and Y, respectively.

Let  $s \in \mathbb{C}$ . If for every  $u_0 \in U$ , there exists a unique solution of the form  $(u_0e^{st}, z_0e^{st}, y_0e^{st})$ , and the mapping  $u_0 \mapsto y_0$  is linear and bounded, then this mapping is called the transfer function at s, and will be denoted by G(s).

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We call a solution of the form  $(u_0e^{st}, z_0e^{st}, y_0e^{st})$  an exponential solution.

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Linear:  $Q(\alpha z_1 + \beta z_2) = \alpha Q z_1 + \beta Q z_2$  for all  $z_1, z_2 \in Z$ ,  $\alpha, \beta \in \mathbb{R}$ , and

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The set of all bounded, linear operators from Z to W is denoted by  $\mathcal{L}(Z,W)$ .

# Transfer function for state linear systems

Consider the abstract differential equation

$$\dot{x}(t) = Ax(t) + Bu(t)$$
  
 $y(t) = Cx(t) + Du(t)$ 

with B,C, and D bounded (linear) operators. Let  $s\in\mathbb{C}$ , and  $u_0\in U$ . We try to find a solution of the form  $(u(t),x(t),y(t))=(u_0e^{st},x_0e^{st},y_0e^{st})$ .

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$$sx_0e^{st} = Ax_0e^{st} + Bu_0e^{st}$$
$$y_0e^{st} = Cx_0e^{st} + Du_0e^{st}.$$

Since  $e^{st}$  is never zero, this is equivalent to:

$$(sI - A)z_0 = Bu_0$$
  
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This clearly defines a bounded linear mapping from  $u_0$  to  $y_0$ , and so the transfer function at s is given by

$$G(s) = C(sI - A)^{-1}B + D.$$

This holds for all  $s \in \rho(A) := \{ s \in \mathbb{C} \mid (sI - A)^{-1} \text{ exists as bounded operator} \}.$ 

## Example

We take a vibrating string with no force at the boundary. We apply a force on it uniformly at one half, and we measure the average position in the other half;

$$\frac{\partial^2 w}{\partial t^2}(\zeta, t) = c^2 \frac{\partial^2 w}{\partial \zeta^2}(\zeta, t) + \mathbb{1}_{\left[\frac{1}{2}L, L\right]}(\zeta)u(t)$$

$$\frac{\partial w}{\partial \zeta}(0, t) = \frac{\partial w}{\partial \zeta}(L, t) = 0$$

$$y(t) = \int_0^{\frac{1}{2}L} w(\zeta, t)d\zeta.$$

To obtain the transfer function, we could follow two approaches.

The p.d.e. can be written as

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This will turn out to be (almost) the same as the equations which need to be solve in method 2.

We try to find an exponential solution of the p.d.e. This gives the following equations

$$s^{2}x_{0}(\zeta)e^{st} = c^{2}\frac{d^{2}x_{0}}{d\zeta^{2}}(\zeta)e^{st} + \mathbb{1}_{\left[\frac{1}{2}L,L\right]}(\zeta)u_{0}e^{st}$$
$$\frac{dx_{0}}{d\zeta}(0)e^{st} = \frac{dx_{0}}{d\zeta}(L)e^{st} = 0$$
$$y_{0}e^{st} = \int_{0}^{\frac{1}{2}L}x_{0}(\zeta)e^{st}d\zeta.$$

Hence

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The first two lines represent an o.d.e. with boundary conditions.

The solution of

$$s^{2}x_{0}(\zeta) = c^{2}\frac{d^{2}x_{0}}{d\zeta^{2}}(\zeta) + \mathbb{1}_{\left[\frac{1}{2}L,L\right]}(\zeta)u_{0}$$
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is given as

$$x_0(\zeta) = \cosh(\frac{s}{c}\zeta)x_0(0) - \frac{1}{sc} \int_0^{\zeta} \sinh(\frac{s}{c}(\zeta - \tau)) \mathbb{1}_{[1/2L,L]}(\tau)u_0 d\tau$$

with

$$x_0(0) = \frac{\sinh(\frac{s}{c}\frac{L}{2})u_0}{s^2\sinh(\frac{s}{c}L)} = \frac{u_0}{2s^2\cosh(\frac{s}{c}\frac{L}{2})}.$$

## Transfer function

Using this we find that

$$y_0 = \int_0^{\frac{1}{2}L} x_0(\zeta) d\zeta$$
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Hence the transfer function is given by

$$G(s) = \frac{c \tanh(\frac{s}{c} \frac{L}{2})}{2s^3}.$$

#### Transfer function, remark

If you write the solution of the o.d.e.

$$s^{2}x_{0}(\zeta) = c^{2}\frac{d^{2}x_{0}}{d\zeta^{2}}(\zeta) + \mathbb{1}_{\left[\frac{1}{2}L,L\right]}(\zeta)u_{0}$$
$$\frac{dx_{0}}{d\zeta}(0) = \frac{dx_{0}}{d\zeta}(L) = 0$$

as a Fourier cosine series, then you find another expression for the transfer function. Namely,

$$G(s) = \frac{L}{4s^2} - 2L \sum_{n=1}^{\infty} \frac{\sin(n\pi\frac{1}{2})^2}{n^2\pi^2(s^2L^2 + n^2\pi^2c^2)}.$$

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However, the transfer function is unique, and so we find that

$$\frac{c \tanh(\frac{s}{c} \frac{L}{2})}{2s^3} = G(s) = \frac{L}{4s^2} - 2L \sum_{n=1}^{\infty} \frac{\sin(n\pi \frac{1}{2})^2}{n^2\pi^2(s^2L^2 + n^2\pi^2c^2)}.$$

#### Transfer functions

So we have seen that working with exponential solutions, directly on the p.d.e., works very well.

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We can do that for systems with control and observation at the boundary.

## Transfer function, boundary control and observation

Example

Consider the system with boundary control and observation

$$\frac{\partial w}{\partial t}(\zeta, t) = \frac{\partial w}{\partial \zeta}(\zeta, t)$$

$$w(1, t) = u(t)$$

$$y(t) = w(0, t).$$

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Substituting exponential functions for all signals, gives

$$sx_0(\zeta)e^{st} = \frac{dx_0}{d\zeta}(\zeta)e^{st}$$

$$x_0(1)e^{st} = u_0e^{st}$$

$$y_0e^{st} = x_0(0)e^{st}.$$

Thus

# Example of transfer function with boundary control and observation

$$sx_0(\zeta) = \frac{dx_0}{d\zeta}(\zeta)$$

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The solution equals  $x_0(\zeta) = e^{s(\zeta-1)}u_0$ . Thus  $y_0 = e^{-s}u_0$ .

The transfer function equals

$$G(s) = e^{-s}$$
  $s \in \mathbb{C}$ .

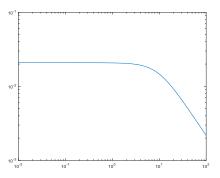
# Bode and Nyquist plots

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$$G(s) = \frac{\tanh(\sqrt{s}/2)}{2s\sqrt{s}} - \frac{1}{4s}$$



Consider the port-Hamiltonian system with input and outputs

$$\frac{\partial x}{\partial t}(\zeta, t) = \left(P_1 \frac{\partial}{\partial \zeta} + P_0\right) \left[\mathcal{H}(\zeta)x(\zeta, t)\right] 
u(t) = W_{B,1} \left[\frac{\mathcal{H}(b)x(b, t)}{\mathcal{H}(a)x(a, t)}\right], \quad 0 = W_{B,2} \left[\frac{\mathcal{H}(b)x(b, t)}{\mathcal{H}(a)x(a, t)}\right], 
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y(t) = W_C \left[\frac{\mathcal{H}(b)x(b,t)}{\mathcal{H}(a)x(a,t)}\right].$$

Assume that the energy balance can be expressed in the inputs and outputs. That is

$$\dot{H}(t) = \left[ u(t)^{\top}, y(t)^{\top} \right] Q \begin{bmatrix} u(t) \\ y(t) \end{bmatrix}$$

with Q a symmetric matrix.

Since exponential solutions are solutions, the power balance

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Remark: Since the s in the exponential solution may be complex, we have to write the power balance for complex valued solutions. The (complex) power balance equals

$$\dot{H}(t) = \left[u(t)^*, y(t)^*\right] Q \begin{bmatrix} u(t) \\ y(t) \end{bmatrix}.$$

Hence for the exponential solution the power balance can be written as

$$\dot{H}(t) = \left[ u(t)^*, y(t)^* \right] Q \begin{bmatrix} u(t) \\ y(t) \end{bmatrix}$$
$$= \left[ u_0^* e^{\overline{s}t}, y_0^* e^{\overline{s}t}, \right] Q \begin{bmatrix} u_0 e^{st} \\ y_0 e^{st} \end{bmatrix}$$

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$$\begin{split} \dot{H}(t) &= \left[u(t)^*, y(t)^*\right] Q \begin{bmatrix} u(t) \\ y(t) \end{bmatrix} \\ &= \left[u_0^* e^{\overline{s}t}, y_0^* e^{\overline{s}t}, \right] Q \begin{bmatrix} u_0 e^{st} \\ y_0 e^{st} \end{bmatrix} \\ &= \left[u_0^*, y_0^*\right] Q \begin{bmatrix} u_0 \\ y_0 \end{bmatrix} e^{2\operatorname{Re}(s)t} \\ &= \left[u_0^*, u_0^* G(s)^*\right] Q \begin{bmatrix} u_0 \\ G(s) u_0 \end{bmatrix} e^{2\operatorname{Re}(s)t}. \end{split}$$

Since

$$H(t) = ||x(t)||_X^2 = \langle x(t), x(t) \rangle_X,$$

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Combining the two results gives that

$$2\operatorname{Re}(s)\|x_0\|_X^2 e^{2\operatorname{Re}(s)t} = \dot{H}(t) = \left[u_0^*, u_0^* G(s)^*\right] Q \begin{bmatrix} u_0 \\ G(s)u_0 \end{bmatrix} e^{2\operatorname{Re}(s)t}.$$

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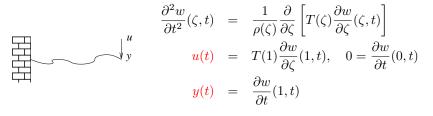
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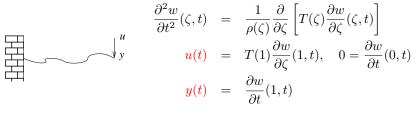
Or equivalently:

$$2\operatorname{Re}(s)\|x_0\|_X^2 = \dot{H}(t) = \left[u_0^*, u_0^*G(s)^*\right] Q \begin{bmatrix} u_0 \\ G(s)u_0 \end{bmatrix}.$$

#### Example: Wave equation

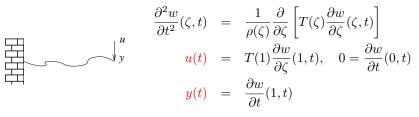


#### Example: Wave equation



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$$\dot{H}(t) = u(t)y(t) = \begin{bmatrix} u(t)^*, y(t)^* \end{bmatrix} \begin{bmatrix} 0 & \frac{1}{2} \\ \frac{1}{2} & 0 \end{bmatrix} \begin{bmatrix} u(t) \\ y(t) \end{bmatrix}.$$

# Transfer function for the vibrating string system

From the general result we find

$$2\operatorname{Re}(s)\|x_0\|_X^2 = \left[u_0^*, u_0^*G(s)^*\right] \begin{bmatrix} 0 & \frac{1}{2} \\ \frac{1}{2} & 0 \end{bmatrix} \begin{bmatrix} u_0 \\ G(s)u_0 \end{bmatrix}$$

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Since  $||x_0||_X^2 \ge 0$  and  $|u_0|^2 > 0$ , we find that for Re(s) > 0 there holds

$$\operatorname{Re}(G(s)) \ge 0$$

Thus G is positive real.